

Quantified Gamow Shell Model interaction for *psd*-shell nuclei

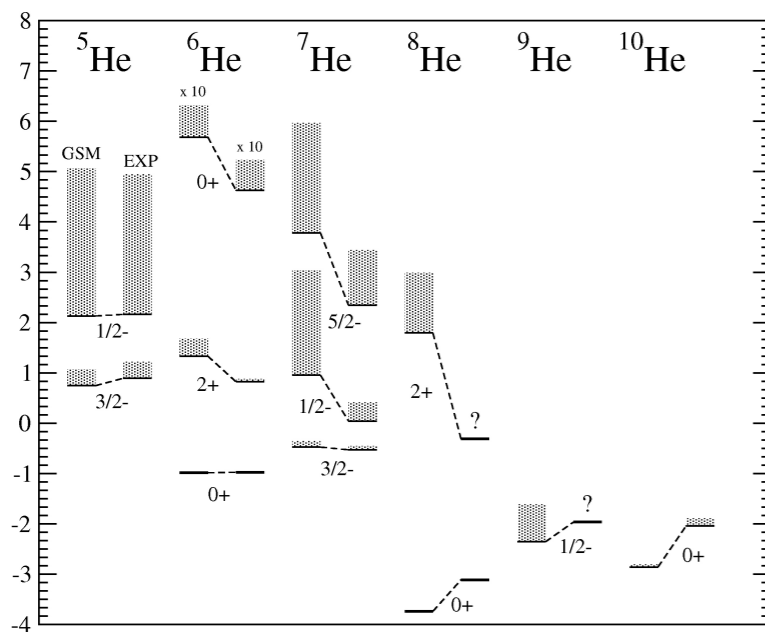
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M. Płoszajczak (GANIL)



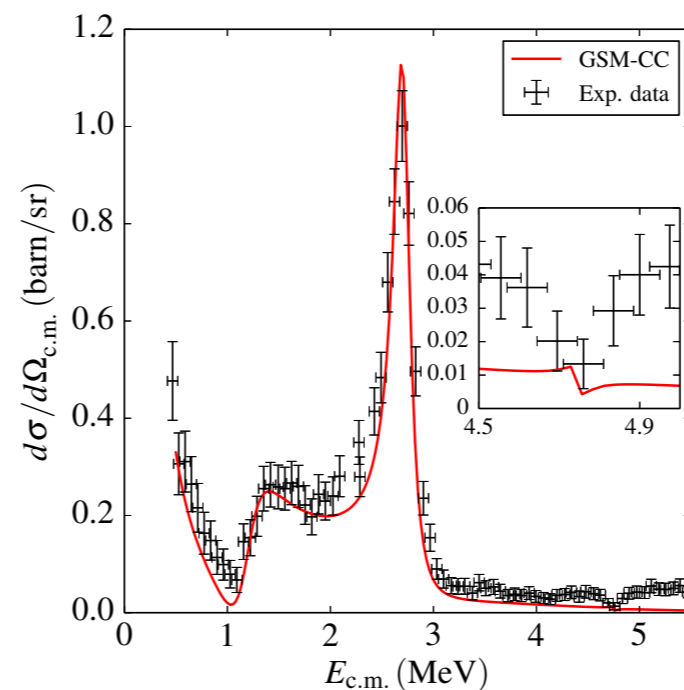
Quantified Gamow Shell Model interaction for *psd*-shell nuclei

We have optimized an **effective N+NN potential** within the **Gamow Shell Model (GSM)** framework, designed to describe a variety of **structure (bound + unbound)** and **reaction** observables across the ***psd*-shell** nuclei ($A \approx 5 - 15$)

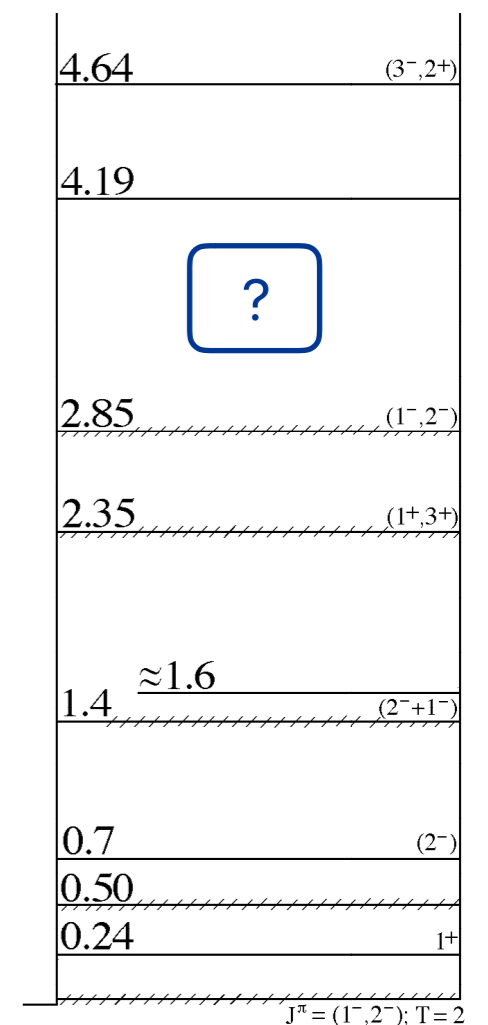
Statistical studies have been carried out to assess **statistical uncertainties and correlations** between parameters and/or predicted observables



GSM spectra of Helium isotopes obtained with a simple interaction



GSM-CC $^{14}\text{O}(p,p')$ excitation function



^{10}Li experimental spectrum

Outline

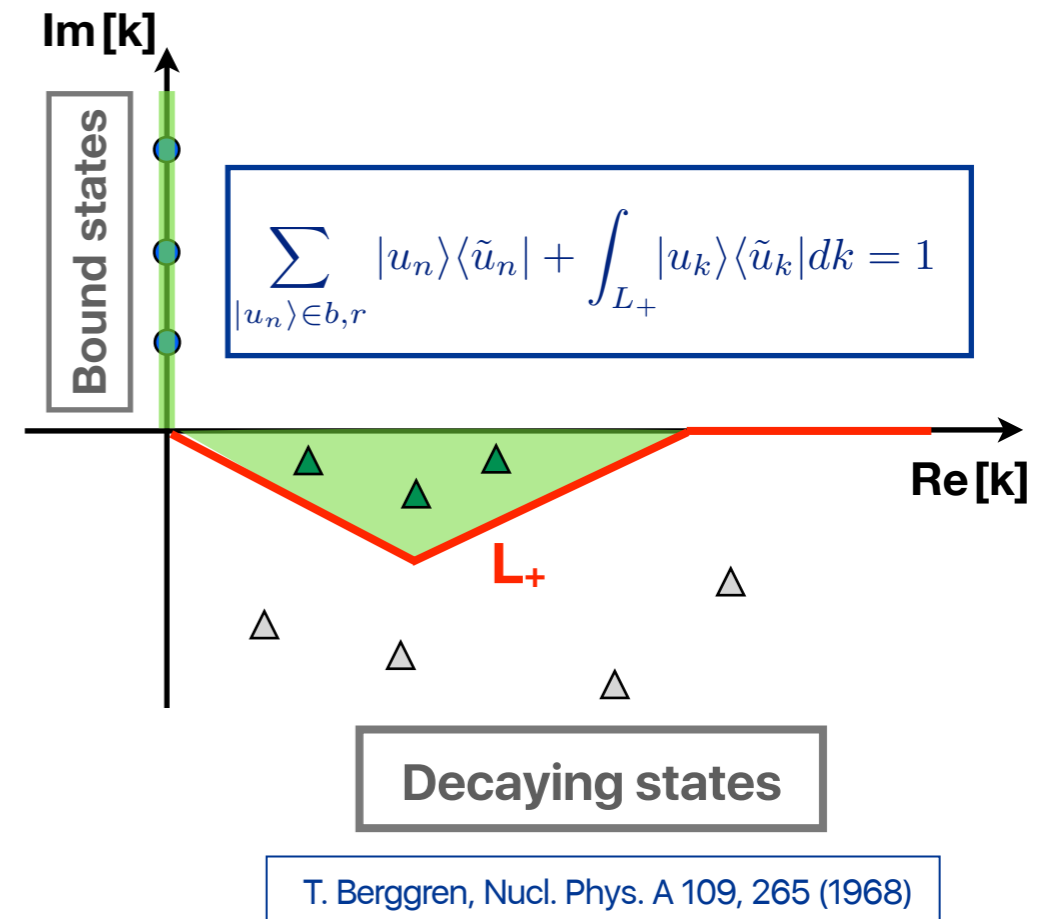
1. The framework
 - The Gamow Shell Model
 - Optimization and Uncertainty Quantification
2. The quantified GSM effective interaction
 - The core potential
 - The NN interaction
3. Applications
 - $A=7$ nuclei
 - Helium spectra
4. Conclusions

- ▶ **Open-quantum system extension** of the traditional Shell Model
- ▶ Both **correlations** and **continuum effects** are treated on the same footing (**Berggren ensemble**)

- ▶ GSM-Cluster orbital shell model (COSM)
Hamiltonian:

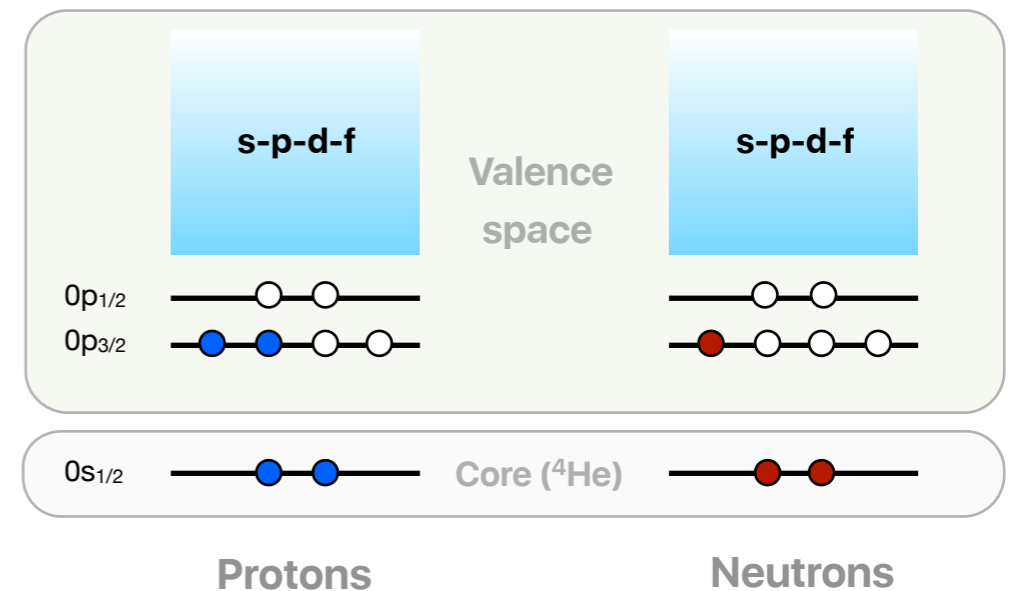
$$H = \sum_{i=1}^{N_v} \left[\frac{\vec{p}_i^2}{2\mu_i} + U(i) \right] + \sum_{i < j=1}^{N_v} \left[V_{res}(i, j) + \frac{\vec{p}_i \cdot \vec{p}_j}{M_c} \right]$$

- ▶ Translational invariance
- ▶ Exact treatment of the Coulomb interaction



The Interaction

- ▶ **^4He core** modeled by a Woods-Saxon + spin-orbit + Coulomb
- ▶ **Configuration space: psdf**
 - $0p_{3/2}$, $0p_{1/2}$ and/or $1s_{1/2}$, $0d_{5/2}$ resonances
 - s, p, d and f scattering continua, $k_{\text{max}} = 2.0 \text{ fm}^{-1}$
- ▶ Effective finite-range NN potential
 - **Gaussian-like with central + spin-orbit + tensor + Coulomb channels**
 - Based on H. Furutani, H. Horiuchi, and R. Tamagaki, Prog. Theor. Phys. 62, 981 (1979)
 - **7 parameters** adjusted to the He, Li, Be chain ground-state energies + chosen excited states
- ▶ The calculations were made possible by the **hybrid parallelization** of the GSM code



The Optimization

- ▶ **Chi-square** minimization

J. Dobaczewski, W. Nazarewicz, P.-G. Reinhard,
J. Phys. G: Nucl. and Part. Phys. 41, 074001 (2014)

$$\chi^2(\mathbf{p}) = \sum_{i=1}^{N_d} \left(\frac{\mathcal{O}_i(\mathbf{p}) - \mathcal{O}_i^{\text{exp}}}{\delta \mathcal{O}_i} \right)^2$$

- ▶ Part of the arbitrariness of the adopted errors is removed by requiring that

$$\frac{\chi^2(\mathbf{p}_0)}{N_{\text{dof}}} \leftrightarrow 1$$

at the minimum \mathbf{p}_0 (similarly to the case of a purely statistical distribution).

- ▶ In the case of a single type of data (and negligible experimental + numerical errors):
⇒ **Global scaling of the adopted errors** (One single Birge factor)

The Optimization

- ▶ The minimization were performed using the **Gauss-Newton** algorithm augmented by the **Singular Value Decomposition (SVD)** technique:

- ▶ **Gauss-Newton method** (for the overdetermined case $N_{\text{data}} > N_{\text{param}}$)

$$P_{(s+1)} = P_{(s)} - \left(J_{(s)}^T J_{(s)} \right)^{-1} J_{(s)}^T F_{(s)}$$

s : step
 P : parameter-vector
 J : Jacobian
 F : residual vector

- ▶ **Singular Value Decomposition (SVD)**: Moore-Penrose pseudoinverse to deal with sloppy parameters:

$$\cancel{M^{-1} = PD^{-1}P^{-1}} \rightarrow M_{MP}^{-1} = PD_{MP}^{-1}P^{-1}$$

$\Rightarrow MX = B$ is solved in a restricted subspace of the space image of M .

The Statistical Uncertainty Quantification

► **Covariance matrix (linear regression)**

$$\mathcal{C} = (J^T J)^{-1}, \quad J_{i\alpha} = \frac{1}{\delta \mathcal{O}_i} \left. \frac{\partial \mathcal{O}_i}{\partial p_\alpha} \right|_{\mathbf{p}_0}$$

Uncertainties
on **parameters**

$$\Delta p_\alpha = \sqrt{\mathcal{C}_{\alpha\alpha}}$$

Uncertainties as in
'reasonable domain'

Uncertainties
on **observables**

$$\overline{\Delta A}^2 = \sum_{\alpha, \beta=1}^{N_p} \left. \frac{\partial A}{\partial p_\alpha} \right|_{\mathbf{p}_0} \mathcal{C}_{\alpha\beta} \left. \frac{\partial A}{\partial p_\beta} \right|_{\mathbf{p}_0}$$

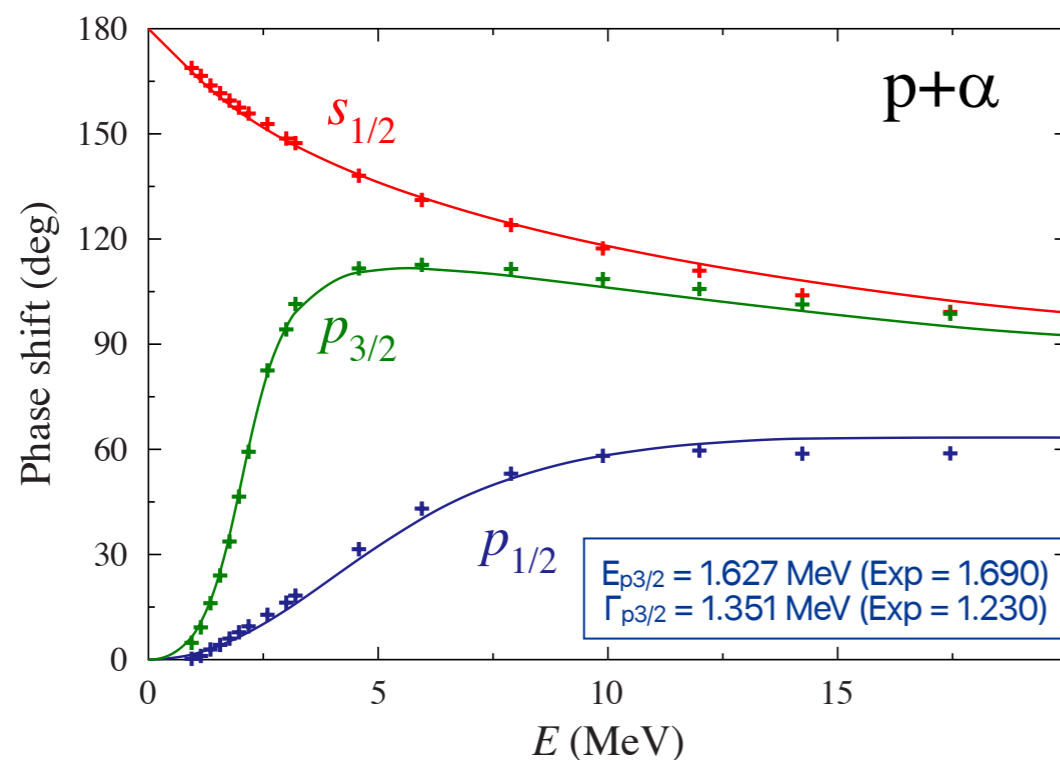
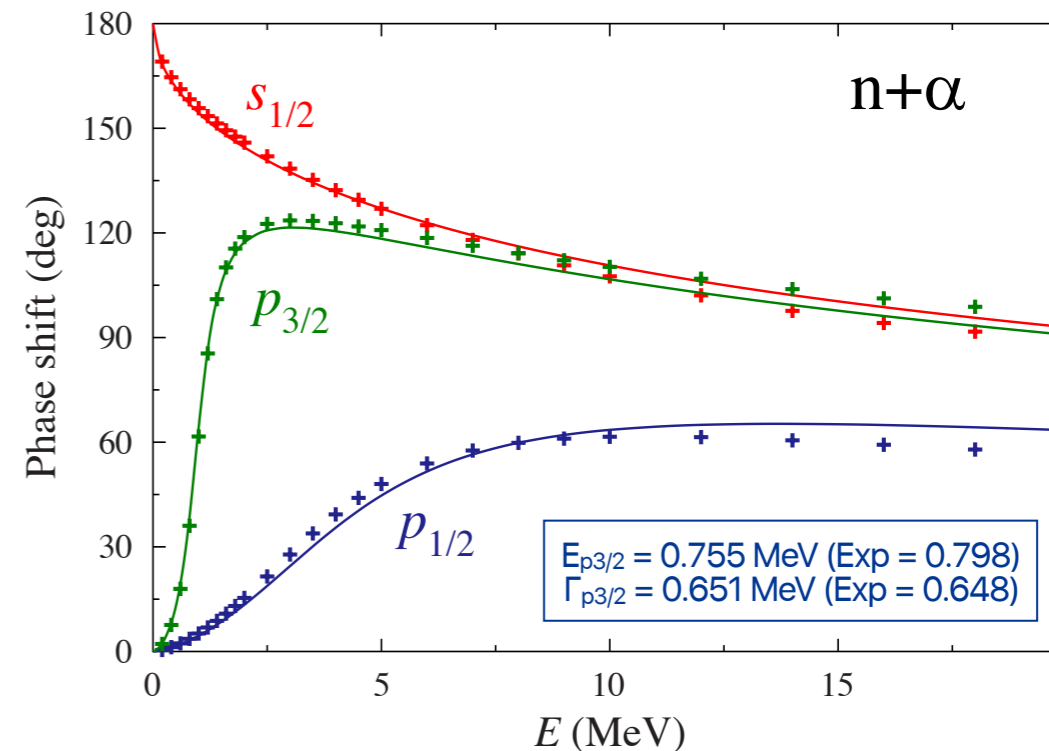
Correlation coefficients
between physical quantities

$$\overline{\Delta A \Delta B} = \sum_{\alpha, \beta=1}^{N_p} \left. \frac{\partial A}{\partial p_\alpha} \right|_{\mathbf{p}_0} \mathcal{C}_{\alpha\beta} \left. \frac{\partial B}{\partial p_\beta} \right|_{\mathbf{p}_0}$$

The Core Potential

- ▶ The Woods-Saxon + spin-orbit + Coulomb was adjusted to N-⁴He phase shifts up to $E_{cm} = 20$ MeV

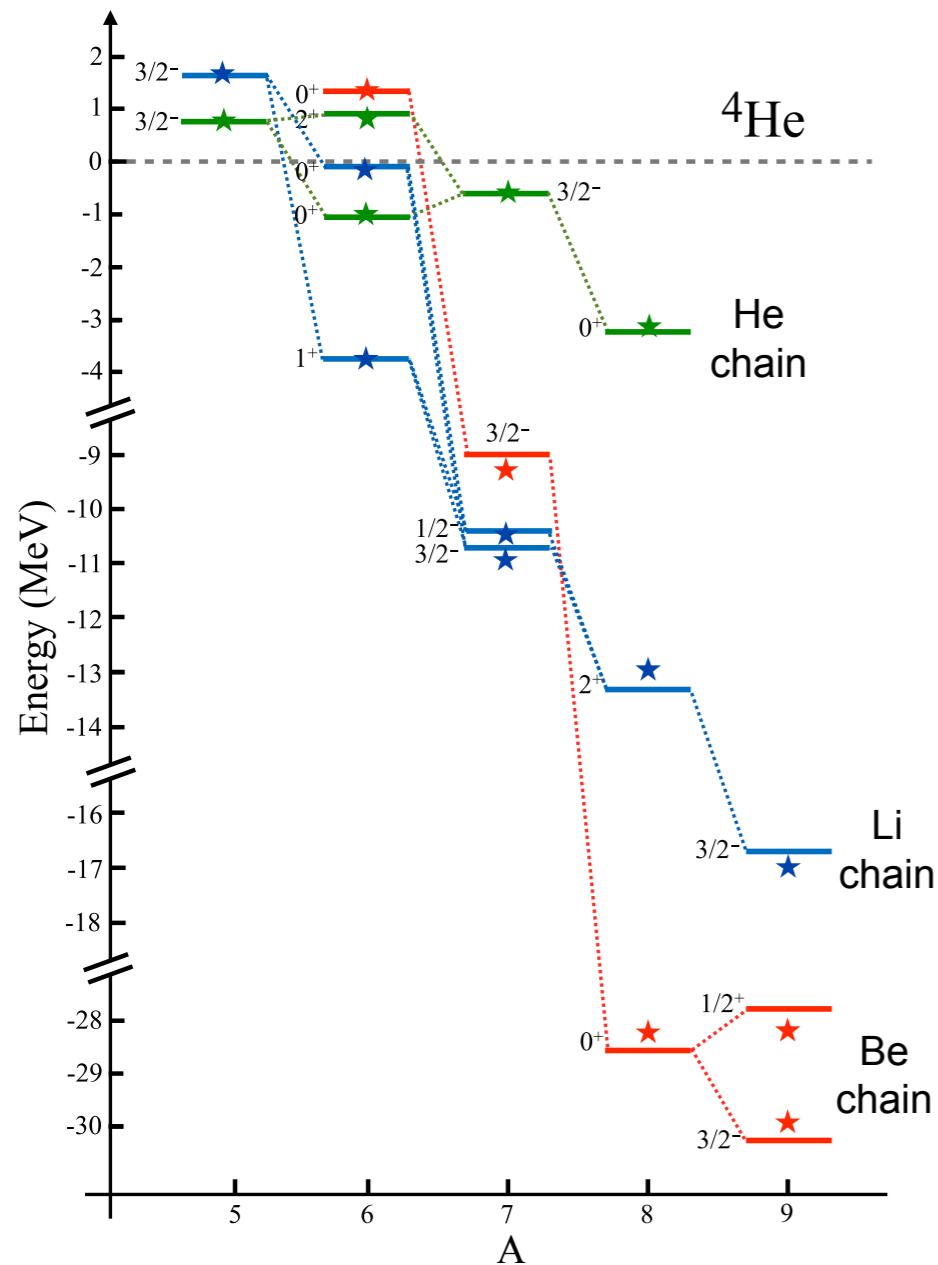
Particle	R_0 [fm]	a [fm]	V_0 [MeV]	V_{so} [MeV fm ²]	R_{ch} [fm]
neutron	2.15 ± 0.04	0.63 ± 0.02	41.9 ± 1.0	7.21 ± 0.20	
proton	2.06 ± 0.04	0.64 ± 0.02	44.4 ± 1.1	7.24 ± 0.21	1.681



n\p	R_0	a	V_0	V_{so}
R_0	•	-0.81	-0.94	-0.78
a	-0.75	•	0.59	0.81
V_0	-0.95	0.51	•	0.62
V_{so}	-0.75	0.84	0.55	•

Correlation coefficients
(Normalized covariance matrix)

The zeroth order NN potential



Nucleus	State	E	E_{exp}	Γ	Γ_{exp}
${}^6\text{He}$	0^+	-1.063	-0.973		
${}^6\text{He}$	2^+	0.938	0.824	168	113(20)
${}^7\text{He}$	$3/2^-$	-0.578	-0.528	178	150(20)
${}^8\text{He}$	0^+	-3.225	-3.112		
${}^6\text{Li}$	1^+	-3.724	-3.699		
${}^6\text{Li}$	0^+	-0.054	-0.136		
${}^7\text{Li}$	$3/2^-$	-10.688	-10.949		
${}^7\text{Li}$	$1/2^-$	-10.359	-10.471		
${}^8\text{Li}$	2^+	-13.350	-12.982		
${}^9\text{Li}$	$3/2^-$	-16.677	-17.046		
${}^6\text{Be}$	0^+	1.390	1.371	21	92(6)
${}^7\text{Be}$	$3/2^-$	-8.977	-9.305		
${}^8\text{Be}$	0^+	-28.572	-28.204	0	0.0056(3)
${}^9\text{Be}$	$3/2^-$	-30.230	-29.870		
${}^9\text{Be}$	$1/2^+$	-27.747	-28.186	0	217(10)

- 4 nucleons in the continuum (converged calculations)
- Could be used with other models (DMRG)

- r.m.s. deviation of 250 keV
- Good starting point for detailed structural and reaction studies

The zeroth order NN potential

Parameters

Parameter		Value
central	S=1, T=1	-3.2 ± 22.0
	S=1, T=0	-5.1 ± 1.0
	S=0, T=0	-21.3 ± 6.6
	S=0, T=1	-5.6 ± 0.5
spin-orbit	S=1, T=1	-540 ± 1240
tensor	S=1, T=1	-12.1 ± 79.5
	S=1, T=0	-14.2 ± 7.1

Singular values (eigenvalues of the normalized Hessian matrix)

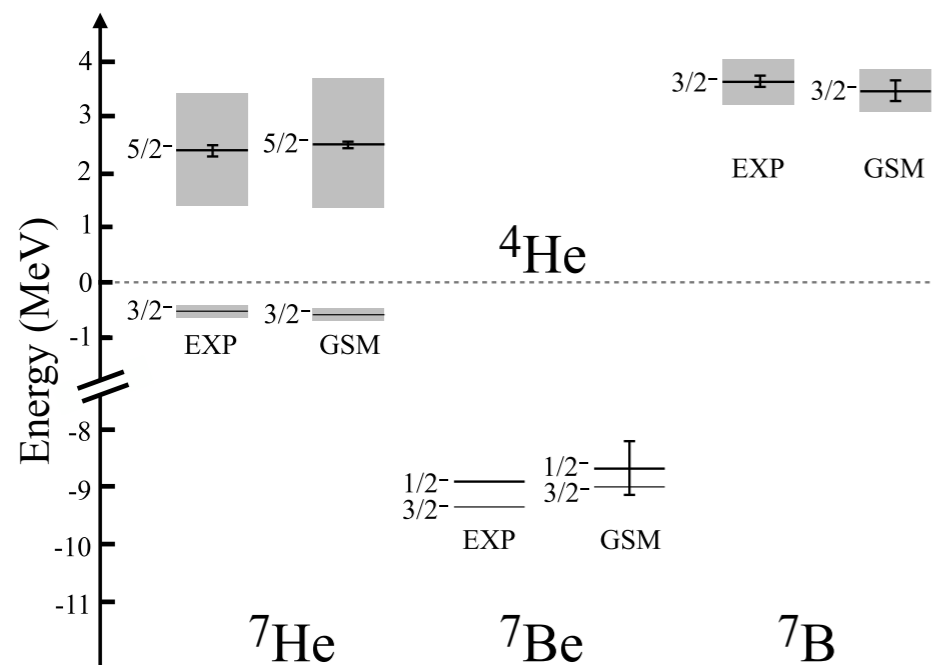
n	s_n	V_c^{11}	V_c^{10}	V_c^{00}	V_c^{01}	V_{LS}^{11}	V_T^{11}	V_T^{10}
1	243	0.00	0.82	-0.03	0.53	0.00	0.00	0.23
2	43.0	0.00	-0.49	-0.02	0.85	0.00	-0.01	-0.19
3	7.06	-0.04	-0.16	0.79	0.05	0.04	-0.07	0.58
4	3.94	0.02	-0.25	-0.61	0.01	-0.09	-0.04	0.75
5	0.57	-0.23	-0.02	-0.09	0.00	0.97	-0.01	0.04
6	0.20	0.65	-0.03	0.04	0.01	0.16	0.74	0.06
7	0.12	0.73	0.01	0.00	0.00	0.16	-0.66	-0.04

- Four parameters completely govern the optimization!
- The three remaining parameters are **sloppy**, i.e. unconstrained by the chosen set of experimental data
 - Can be constrained by experimental data of different kinds (charge/matter radii, EM moments)
 - Could be used to locally fine-tune the interaction

Predictions - Energy Spectra, A=7 nuclei

- ▶ Computed uncertainties have **two components**: $\Delta E = \sqrt{\Delta E_N^2 + \Delta E_{NN}^2}$

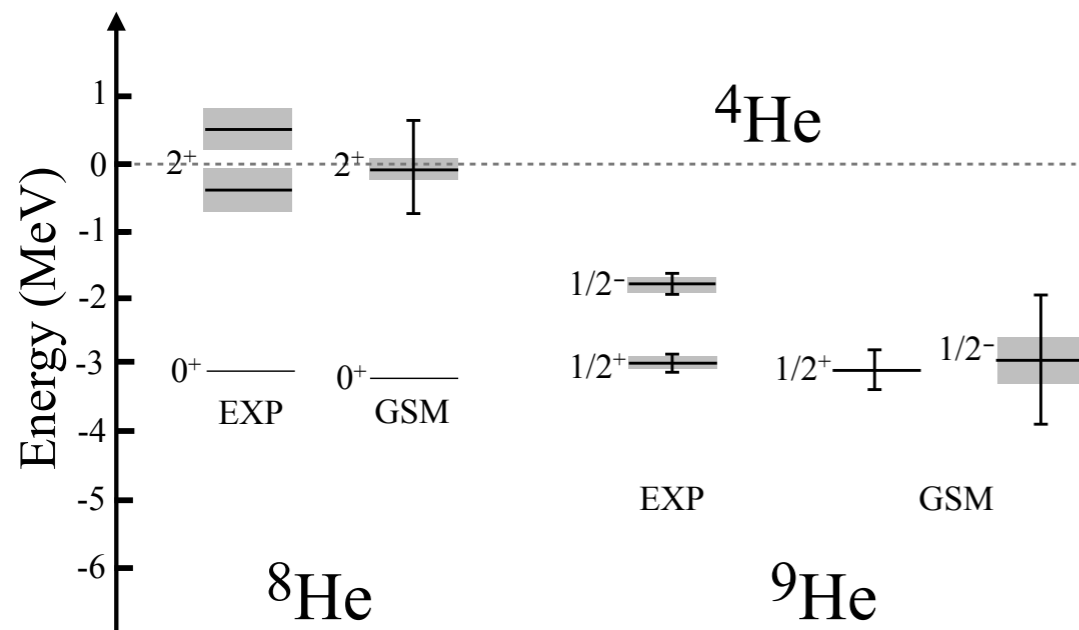
- ▶ **A=7 nuclei:**



State	E_{calc} (MeV)	E_{exp} (MeV)	Γ_{calc} (keV)	Γ_{exp} (keV)
${}^7\text{He}, 5/2^-$	+2.50 (2)	+2.39 (9)	2250 (280)	1990 (170)
${}^7\text{Be}, 1/2^-$	-8.67 (45)	-8.88		
${}^7\text{B}, 3/2^-$	+3.42 (21)	+3.58 (7)	740 (450)	801 (20)

- ▶ $\Delta E_N < 0.07$ MeV (small compared to ΔE_{NN})
- ▶ **Good overall agreement** for the energies and the widths

Predictions - Energy Spectra, Heliums



State	E_{calc} (MeV)	E_{exp} (MeV)	Γ_{calc} (keV)	Γ_{exp} (keV)
$^8\text{He}, 2^+$	-0.10 (75)	-0.41 / +0.49	290 (1010)	1990 (170)
$^9\text{He}, 1/2^+$	-3.12 (31)	-2.93 (9)*	0	180 (160)*
$^9\text{He}, 1/2^-$	-2.98 (102)	-1.88 (12)*	630 (330)	130 (170)*

[*] Al Kalanee et al., PRC 88, 034301 (2013)

▶ $^8\text{He}, 2^+$:

- Uncertainty coming from the WS: $\Delta E_N = 0.07$ MeV
- Prediction **with uncertainties** does not favor any of the experimental scenario

▶ ^9He : not resolved experimentally

- $\Delta E_N (1/2^+) = 0.13$ MeV, $\Delta E_N (1/2^-) = 0.59$ MeV (s1/2 and p1/2 are less constrained in the WS)
- **Shell inversion** observed (mean values)

Summary and Outlook

- ▶ The interaction is well optimized to the ground state + some excited state energies of the He, Li and Be chains
- ▶ The SVD analysis allowed to pinpoint the four interaction parameters which are **reasonably constrained** by the binding energies.
- ▶ The remaining three parameters are **sloppy**; hence new data are needed to limit them.
- ▶ With the covariance matrices calculated for the one and two-body potentials, we assessed **uncertainties** and **correlations** between physical quantities.
- ▶ In collaboration with the Department of Statistics and Probability at MSU, we have initiated a project to perform a **Bayesian study** of the interaction for a fully consistent statistical analysis (see Léo Neufcourt's talk this afternoon for a preliminary Bayesian study of the interaction).

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Thank you for your attention!



Back-up

$$V = V_c + V_{LS} + V_T + V_{\text{Coul}}.$$

$$\tilde{V}_c(r) = \sum_{n=1}^3 V_c^n (W_c^n + B_c^n P_\sigma - H_c^n P_\tau - M_c^n P_\sigma P_\tau) e^{-\beta_c^n r^2} \quad (5)$$

$$\tilde{V}_{LS}(r) = \mathbf{L} \cdot \mathbf{S} \sum_{n=1}^2 V_{LS}^n (W_{LS}^n - H_{LS}^n P_\tau) e^{-\beta_{LS}^n r^2} \quad (6)$$

$$\tilde{V}_T(r) = S_{ij} \sum_{n=1}^3 V_T^n (W_T^n - H_T^n P_\tau) r^2 e^{-\beta_T^n r^2}, \quad (7)$$

where $r \equiv r_{ij}$ stands for the distance between the nucleons i and j , \mathbf{L} is the relative orbital angular momentum, $\mathbf{S} = (\boldsymbol{\sigma}_i + \boldsymbol{\sigma}_j)/2$, $S_{ij} = 3(\boldsymbol{\sigma}_i \cdot \hat{r})(\boldsymbol{\sigma}_j \cdot \hat{r}) - \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j$, and P_σ and P_τ are spin and isospin exchange operators, respec-

In order to be applied in the present GSM formalism, the interaction is rewritten in terms of the spin-isospin projectors Π_{ST} [51, 52]:

$$V_c(r) = V_c^{11} f_c^{11}(r) \Pi_{11} + V_c^{10} f_c^{10}(r) \Pi_{10} + V_c^{00} f_c^{00}(r) \Pi_{00} + V_c^{01} f_c^{01}(r) \Pi_{01}, \quad (8)$$

$$V_{LS}(r) = (\mathbf{L} \cdot \mathbf{S}) V_{LS}^{11} f_{LS}^{11}(r) \Pi_{11}, \quad (9)$$

$$V_T(r) = S_{ij} [V_T^{11} f_T^{11}(r) \Pi_{11} + V_T^{10} f_T^{10}(r) \Pi_{10}], \quad (10)$$