

# THE DATA PERSPECTIVE ON CHIRAL EFFECTIVE FIELD THEORY

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Sweden

ISNET-5 workshop, York,  
Nov. 6-9, 2017

# MANY THANKS TO MY COLLABORATORS

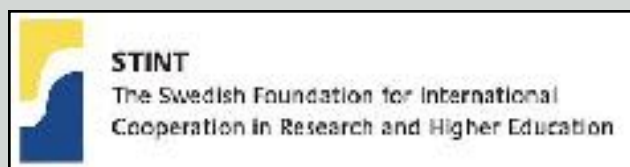
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- ❖ **Boris Carlsson, Andreas Ekström** (Chalmers)
- ❖ Kai Hebel (Darmstadt), Gustav Jansen (ORNL), Kyle Wendt (ORNL/UT now LLNL)

And many people in the *ab initio nuclear theory* community for enlightening discussions

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- STINT
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# INTRODUCTION

# ONE-BULLET OVERVIEW

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- ▶ What information does available nuclear scattering data<sup>1,2</sup> impose<sup>3</sup> on a state-of-the-art “model”<sup>4</sup> of the strong force between nucleons?

<sup>1</sup> only  $\pi N$  and  $NN$

<sup>2</sup> plus  $A=2-3$  bound-state data

<sup>3</sup> using a frequentist approach.

<sup>4</sup> chiral effective field theory => systematically improvable

# FIVE-BULLET OVERVIEW (FOR PHYSICISTS)

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- ▶ Optimize LECs at N3LO of chiral EFT (Weinberg PC) with respect to  $\pi N$ ,  $NN$  scattering data and  $A=2-3$  bound states.
- ▶ There are many LECs – making this a very difficult optimization problem.
- ▶ We attempt uncertainty quantification and error propagation to few-body observables.
- ▶ The analysis is repeated using information from a Roy-Steiner analysis of  $\pi N$  data with resulting constraints on  $c_i:s$ ,  $d_i:s$ ,  $e_i:s$ .
- ▶ Summary and Outlook

# FIVE-BULLET OVERVIEW (TRANSLATED)

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- ▶ Optimize LECs at N3LO of chiral EFT (Weinberg PC) with respect to  $\pi N$ ,  $NN$  scattering data and  $A=2-3$  bound states.

**Translation:** 41 model parameters. 6000 experimental data points with small-to-large error bars. There is some prior knowledge of the model error.

- ▶ There are many LECs – making this a very difficult optimization problem.

**Translation:** We encounter a very flat  $\chi^2$ -surface with many local minima (sometimes non-quadratic).

**Question:** How could this optimization problem be handled? (see also Andreas' talk)

# FIVE-BULLET OVERVIEW (TRANSLATED)

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- ▶ We attempt uncertainty quantification and error propagation to few-body observables.

**Question:** How to best perform uncertainty quantification and error propagation in this situation?

- ▶ The analysis is repeated using information from a Roy-Steiner analysis of  $\pi N$  data with resulting constraints on  $c_i:s$ ,  $d_i:s$ ,  $e_i:s$ .

**Question:** How to handle the errors when combining experimental data with information from a theory analysis?

- ▶ Summary and Outlook

**Translation:** Summary and Outlook

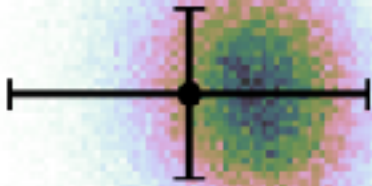
## NEW ARTICLE

## Uncertainty Analysis and Order-by-Order Optimization of Chiral Nuclear Interactions

B.D. Carlsson *et al.*

Phys. Rev. X 6, 011019 (2016)

Estimating errors is a fundamental component of science. Researchers present a new framework for quantifying uncertainties associated with nuclear interactions for low-mass nuclei.



Based on work presented in:

- B. Carlsson, A. Ekström, **CF** et al., Phys. Rev. X 6 (2016) 011019
- B. Carlsson, PhD thesis, manuscript in preparation

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**FROM  $\pi$ N AND NN TO A=4 WITH CHIRAL  
EFT AND ERROR ANALYSIS**

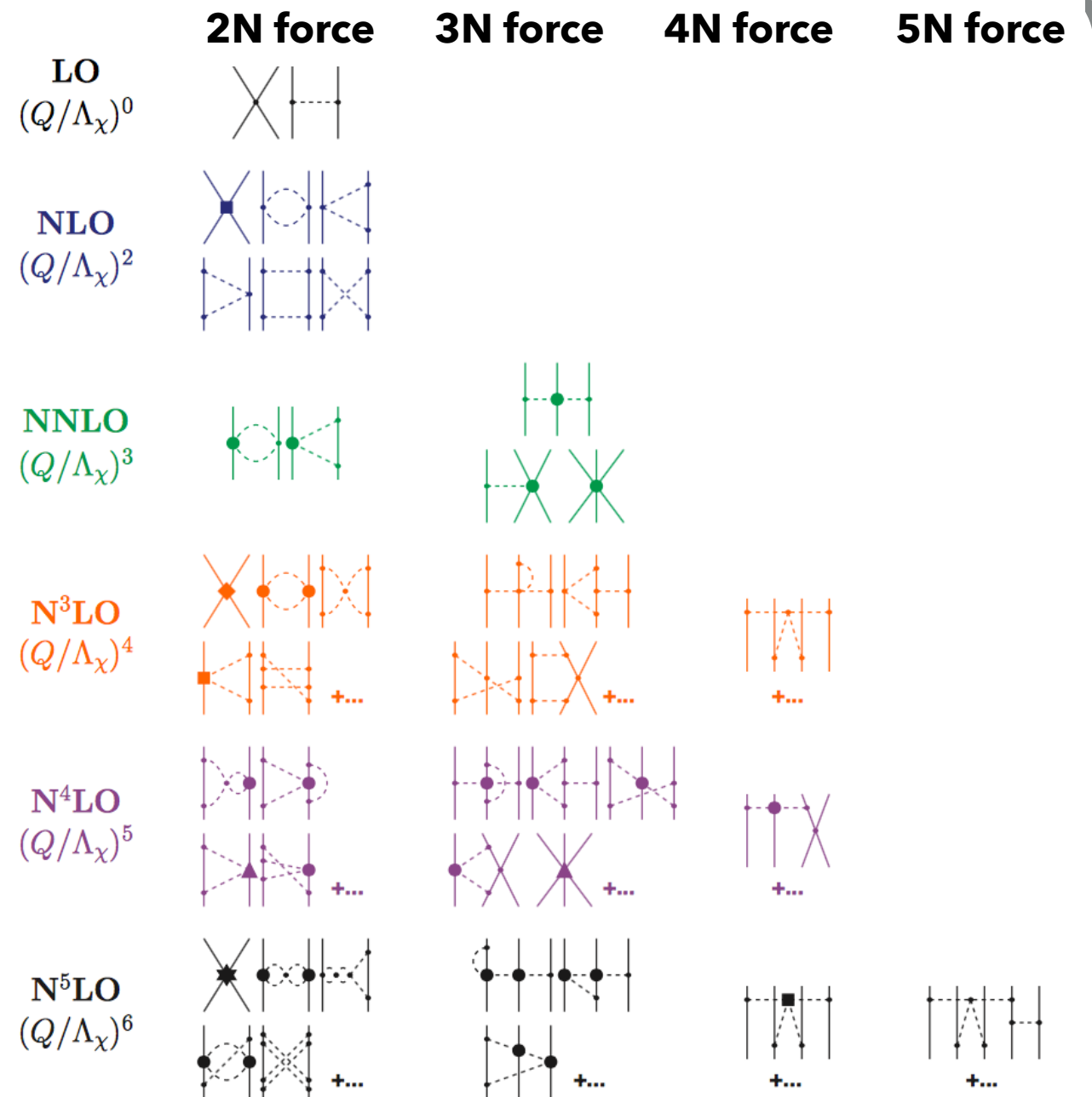


# CHIRAL EFT FOR NUCLEAR INTERACTIONS

## Chiral EFT

- Systematic low-energy expansion:  $(q/\Lambda_\chi)^\nu$
- Connects several sectors:  $\pi N$ ,  $NN$ ,  $NNN$ ,  $j_N$
- (Unknown) short-range physics included as contact interactions.
- LECs need to be fitted to data.

$$\chi^2(\vec{p}) = \sum_i \left( \frac{O_i^{\text{theo}}(\vec{p}) - O_i^{\text{exp}}}{\sigma_{\text{tot},i}} \right)^2$$

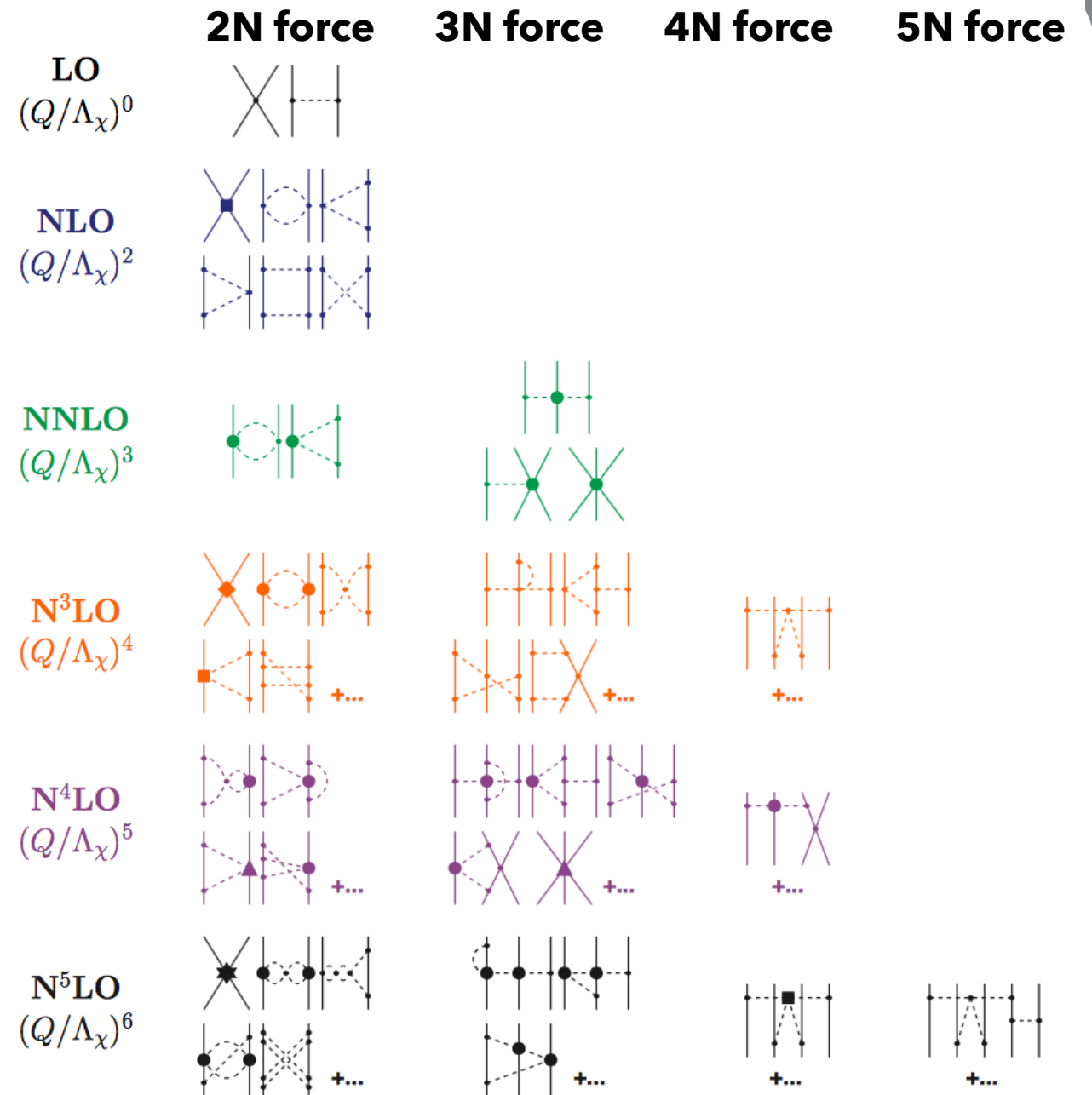


See Evgeny's presentation

# CHIRAL EFT FOR NUCLEAR INTERACTIONS

## Selected key results

- Non-local, 2NF up to N5LO
  - D. R. Entem et al. PRC 91, 014002 (2015)
  - R. Machleidt et al. Phys. Rep. 503 (2011)
- Non-local, leading 3NF
  - E. Epelbaum et al. PRC 66, 064001 (2002)
  - K. Hebeler et al. PRC 91, 044001 (2015)
- Non-local, sub-leading 3NF
  - V. Bernard et al. PRC 77, 064004 (2008)
  - V. Bernard et al. PRC 84, 054001 (2011)
  - K. Hebeler et al. PRC 91, 044001 (2015)
- 4th-order  $\pi N$  scattering
  - H. Krebs et al. PRC 85, 054006 (2015)



See Evgeny's presentation

# CHIRAL NUCLEAR INTERACTIONS

## Number of LECs

- Non-local, 2NF up to N3LO

**26 contacts**

**4  $c_{i:s}$  + 4  $d_{i:s}$**

- Non-local, leading 3NF

**$C_D, C_E$**

3  $c_{i:s}$

- Non-local, subleading 3NF

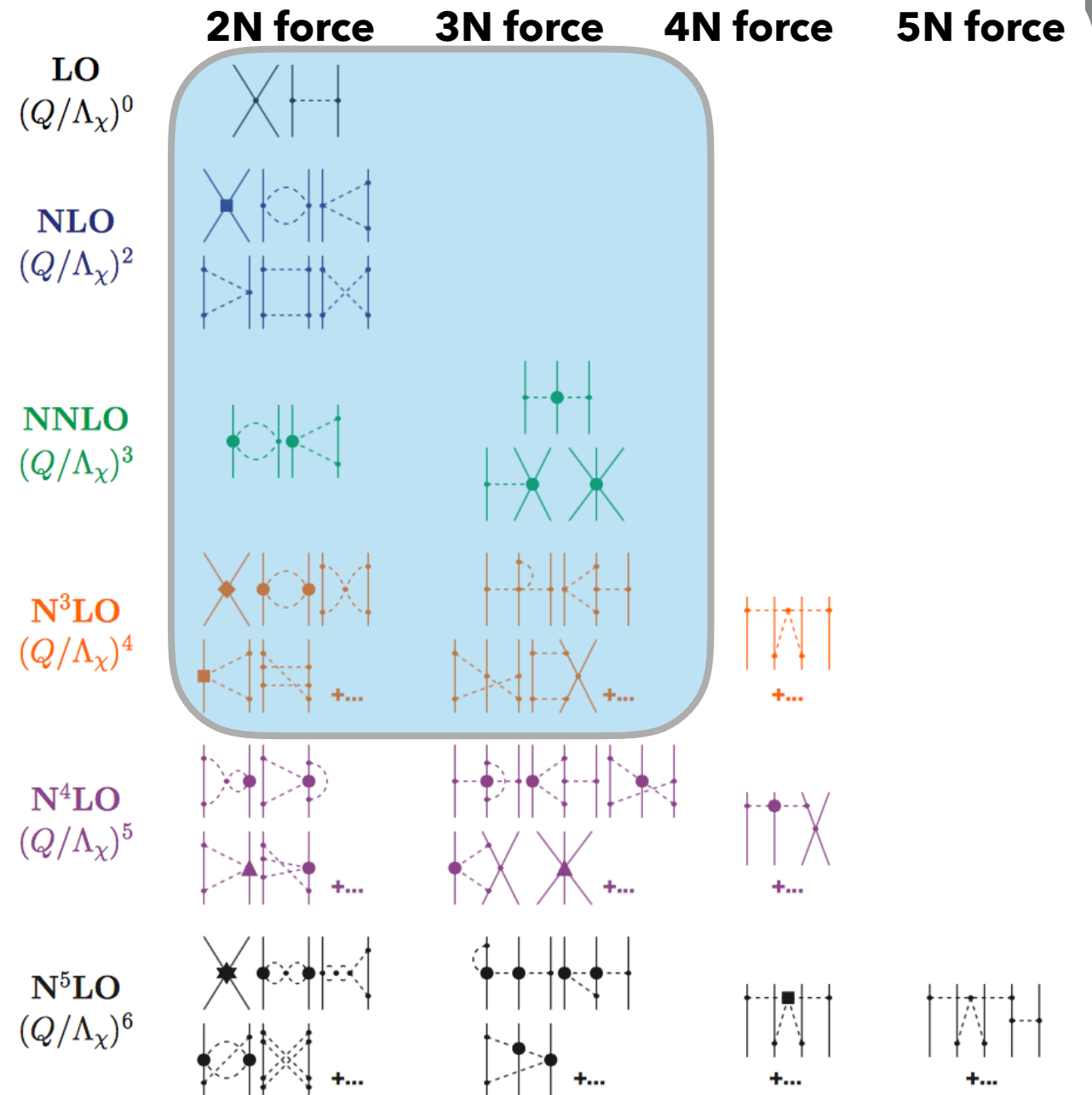
$C_S, C_T$

- 4th-order  $\pi N$  scattering

$c_{i:s} + d_{i:s}$

**5  $e_{i:s}$**

**TOTAL:  
41 parameters**



# INPUT AND TECHNOLOGY

## $\pi$ N scattering

- WI08 database
- $T_{\text{lab}}$  between 10-70 MeV
- $N_{\text{data}} = 1347$
- R. Workman et al. (2012)

## NN scattering

- Granada '13 database
- $T_{\text{lab}}$  between 0-290 MeV
- $N_{\text{data}} = 4753$  (np + pp)
- R. Navarro Pérez et al. (2013)

**All 6000 residuals computed on 1 node in ~90 sec.**

## **A=2,3 bound states**

- ${}^2\text{H}, {}^3\text{H}, {}^3\text{He}$  [binding energy, radius,  $Q({}^2\text{H}), {}^3\text{H}$  half life]

**On 1 node in ~10 sec**

**+ derivatives! ( $\times 2-20$  cost)**

**Alternatively... theoretical analysis of data**

## $\pi$ N scattering

- Roy-Steiner analysis  
M. Hoferichter et al. (2015)

## NN scattering

- Phase shifts from partial wave analysis

# OPTIMIZATION STRATEGY

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**Low-energy constants (LECs) need to be fitted to experimental data.**

$$\chi^2(\vec{p}) \equiv \sum_i \left( \frac{O_i^{\text{theo}}(\vec{p}) - O_i^{\text{expr}}}{\sigma_{\text{tot},i}} \right)^2 \equiv \sum_i r_i^2(\vec{p})$$

- ▶ Derivative-free optimization using POUNDerS was used in our earliest works
- ▶ More efficient minimization algorithms (Levenberg-Marquardt, Newton), and statistical error analysis require **derivatives**

$$\frac{\partial r_i}{\partial p_j} \quad \text{and} \quad \frac{\partial^2 r_i}{\partial p_j \partial p_k}$$

- ▶ Numerical derivation using **finite differences** is plagued by **low numerical precision** and is **computationally costly**.
- ▶ Instead, we use **Automatic Differentiation (AD)**

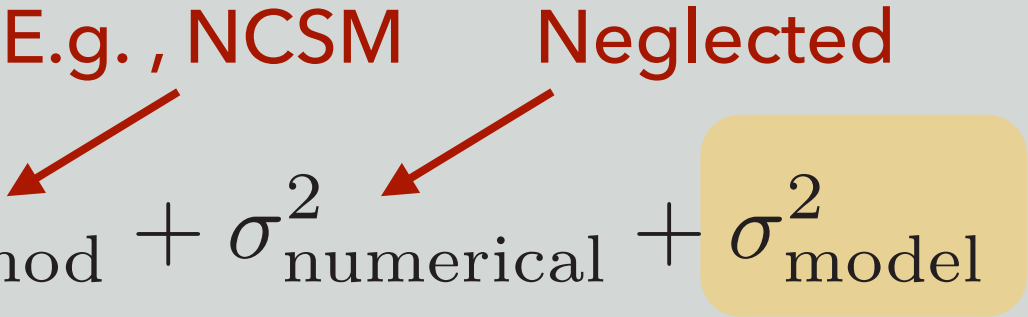
# TOTAL ERROR BUDGET

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- ▶ The total error budget is

$$\sigma_{\text{tot}}^2 = \sigma_{\text{exp}}^2 + \sigma_{\text{method}}^2 + \sigma_{\text{numerical}}^2 + \sigma_{\text{model}}^2$$

E.g., NCSM      Neglected



- ▶ At a given chiral order  $\nu$ , the omitted diagrams should be of order

$$\mathcal{O}\left(\left(Q/\Lambda_\chi\right)^{\nu+1}\right)$$

- ▶ Still needs to be converted to actual numbers  $\sigma_{\text{model}}$
- ▶ We translate this EFT knowledge into an error in the scattering amplitudes

$$\sigma_{\text{model},x}^{(\text{amp})} = C_x \left(\frac{Q}{\Lambda_\chi}\right)^{\nu+1}, \quad x \in \{NN, \pi N\}$$

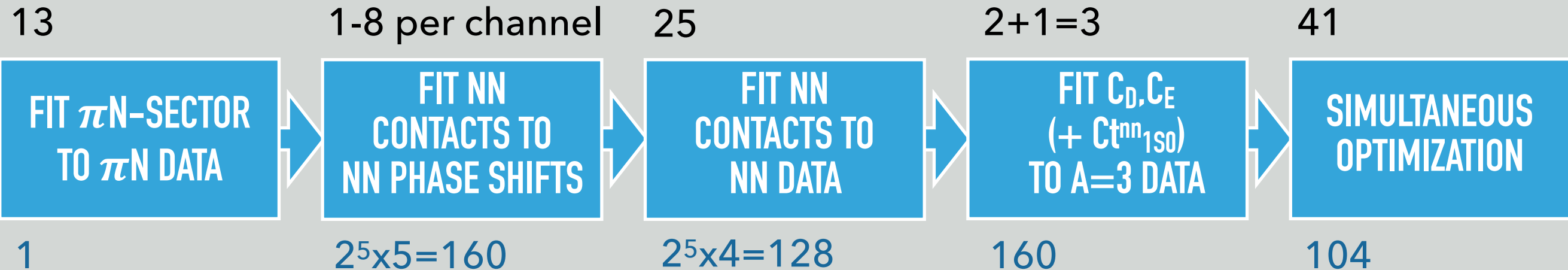
- ▶ which is then propagated to an error in the observable.

# OPTIMIZATION STRATEGY

Low-energy constants (LECs) need to be fitted to experimental data.

$$\chi^2(\vec{p}) \equiv \sum_i r_i^2(\vec{p}) = \sum_{j \in NN} r_j^2(\vec{p}) + \sum_{k \in \pi N} r_k^2(\vec{p}) + \sum_{l \in 3N} r_l^2(\vec{p})$$

# parameters that are allowed to vary:



# minima at each stage

(used as starting points for the next stage); ( $\Lambda=500$  MeV)

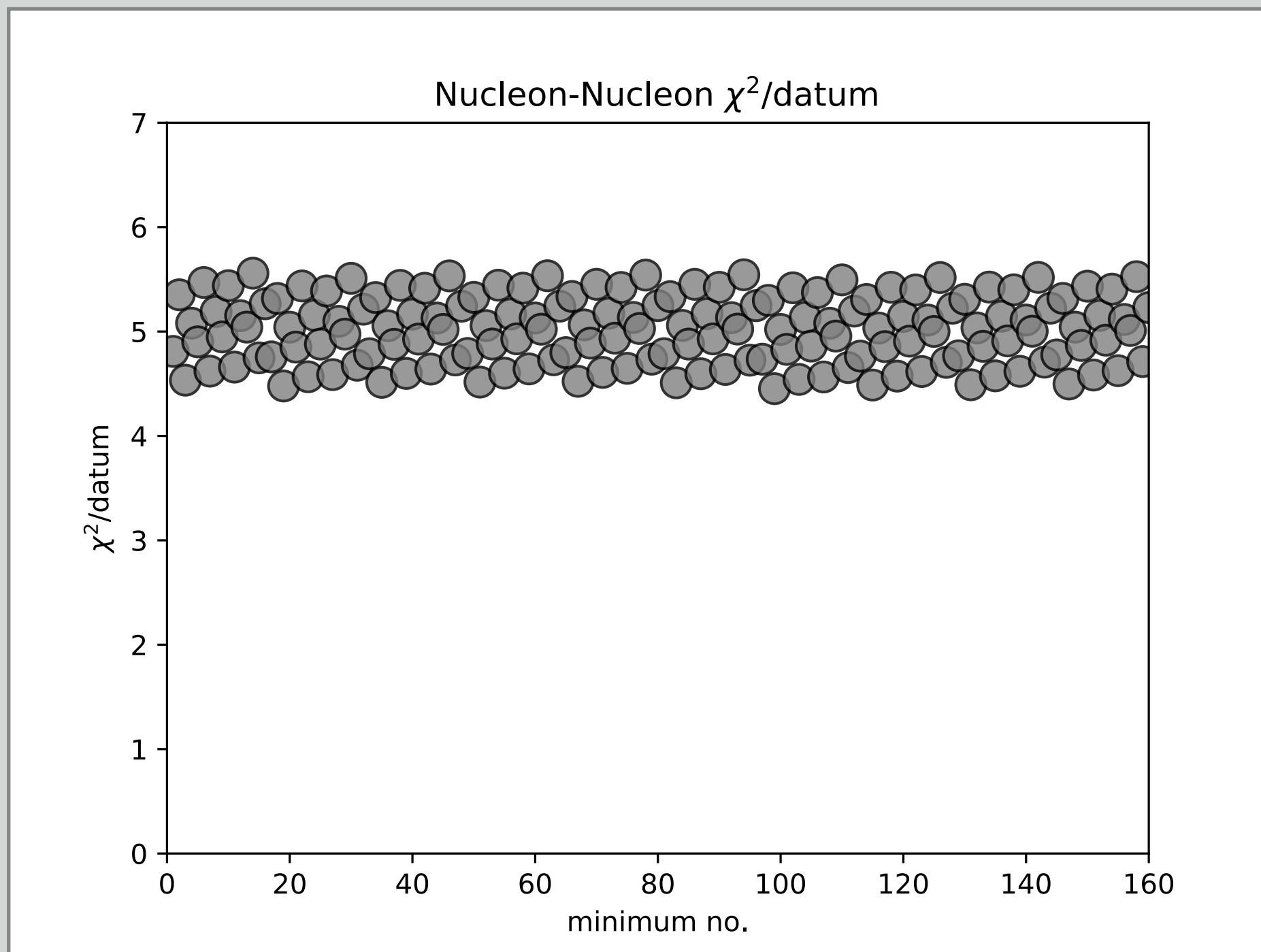
$c_1$	-0.69(50)
$c_2$	+3.0(14)
$c_3$	-4.12(32)
$c_4$	+5.35(81)
$d_1+d_2$	+6.22(44)
$d_3$	-5.31(30)
$d_5$	-0.46(18)
$d_{14}-d_{15}$	-11.00(42)
$e_{14}$	-0.63(95)
$e_{15}$	-7.7(26)
$e_{16}$	+5.9(49)
$e_{17}$	+2.1(18)
$e_{18}$	-8.1(42)

2 per channel:  $^1S_0, ^1P_1, ^3P_0, ^3P_1, ^3P_2-^3F_2$

5 in the deuteron channel:  $^3S_1, ^3D_1, ^3S_1-^3D_1$

2 per channel: Some cases with 2  $c_D, c_E$  optima

4 in the deuteron channel:





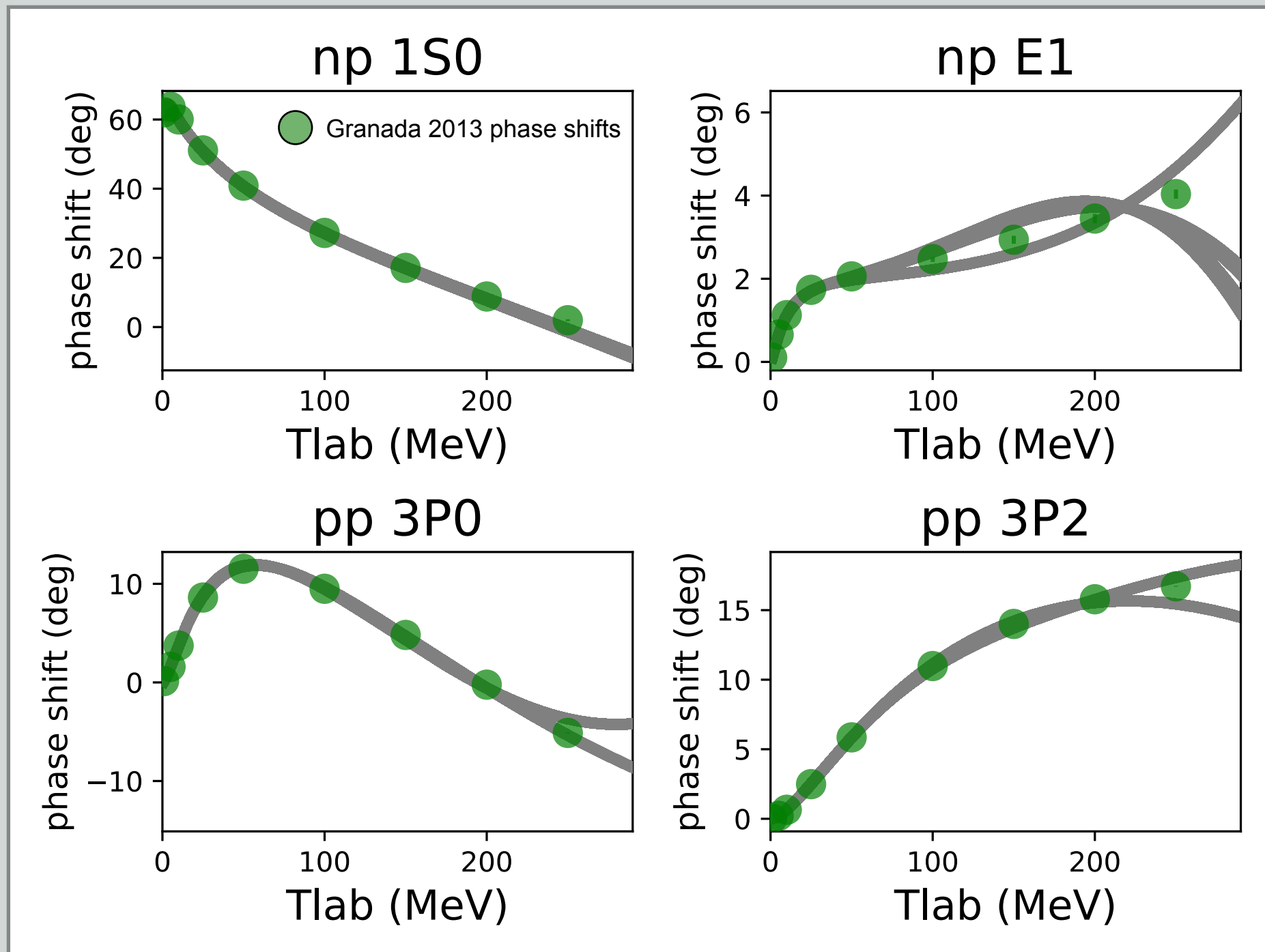
FIT  $\pi$ N-SECTOR  
TO  $\pi$ N DATA

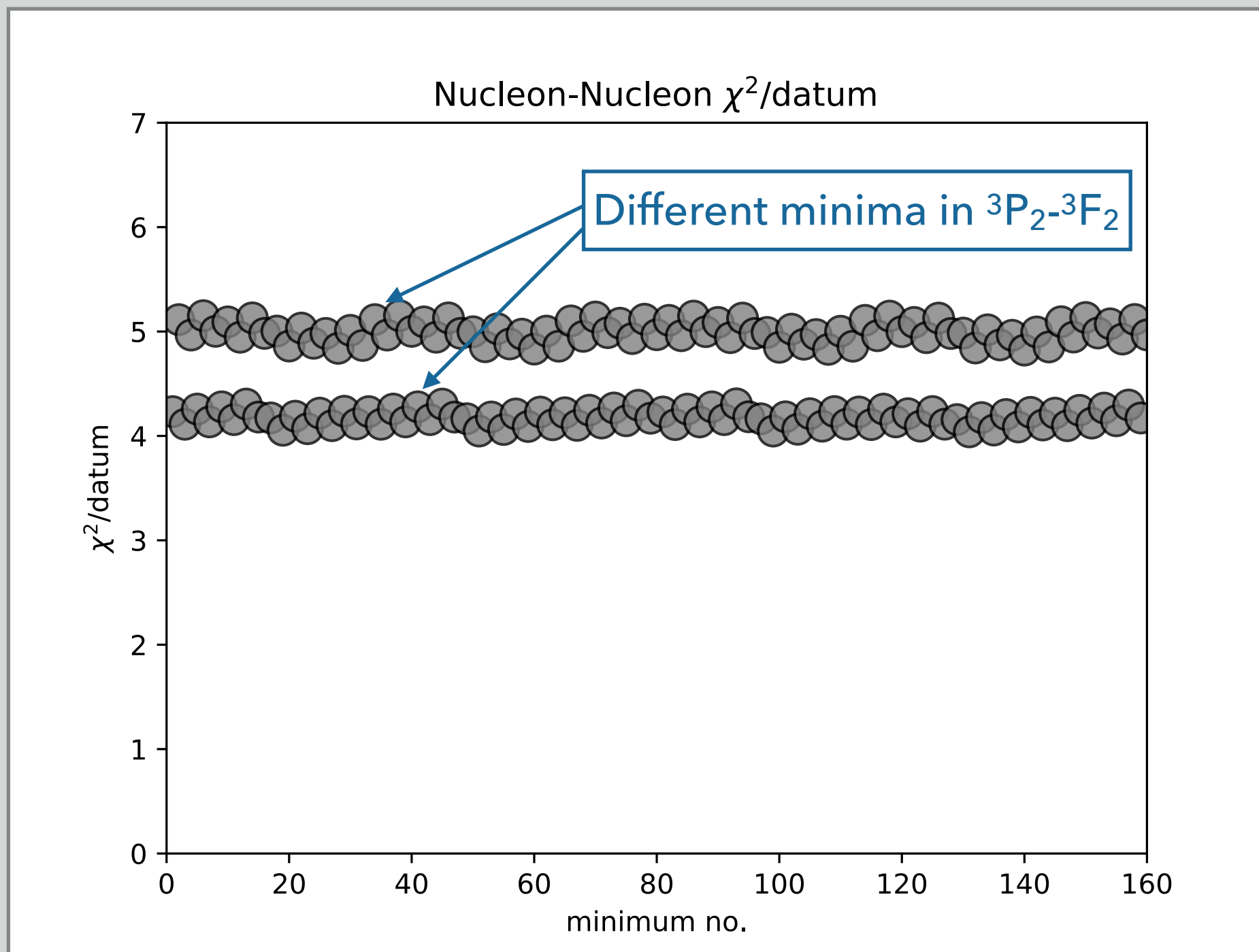
FIT NN  
CONTACTS TO  
NN PHASE SHIFTS

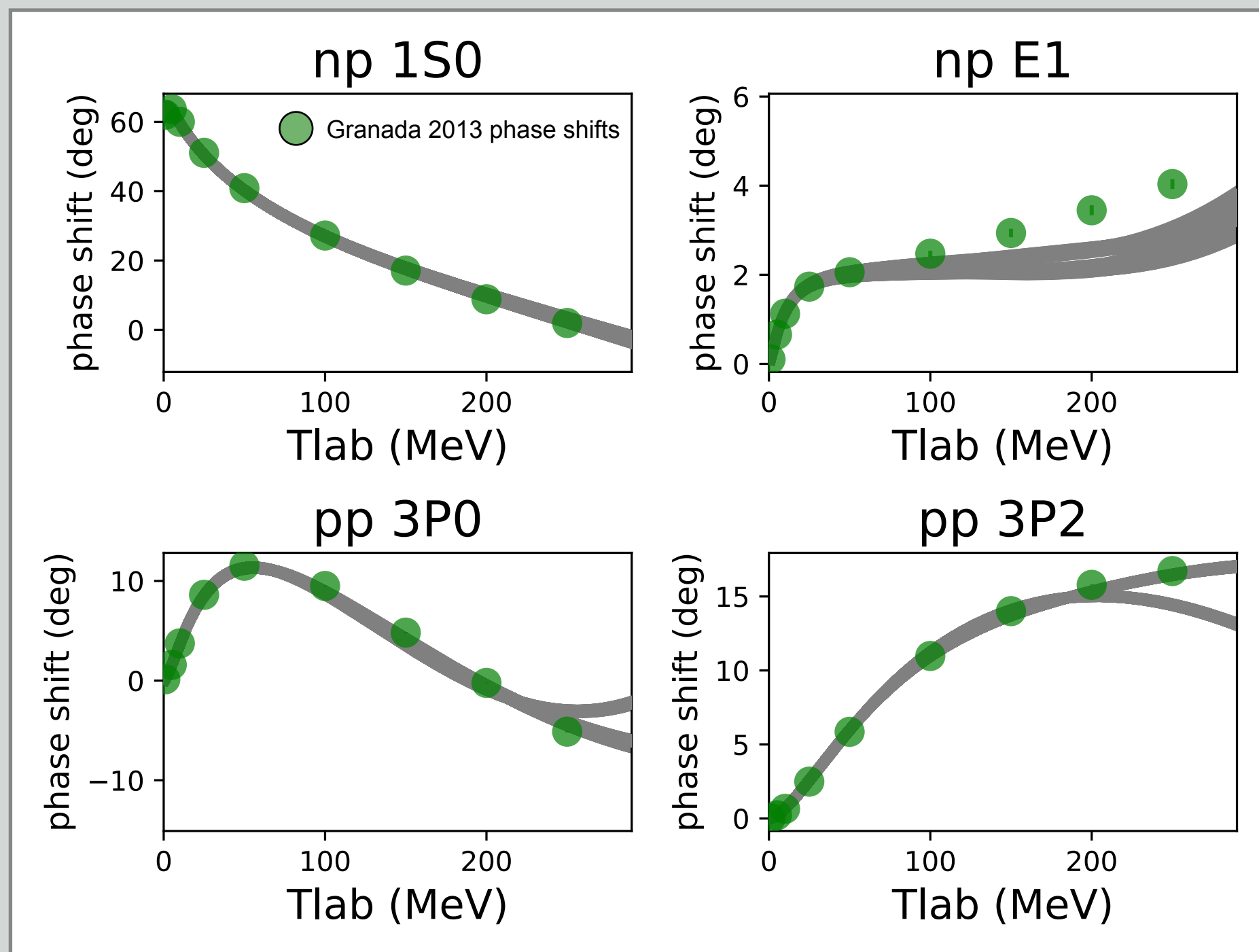
FIT NN  
CONTACTS TO  
NN DATA

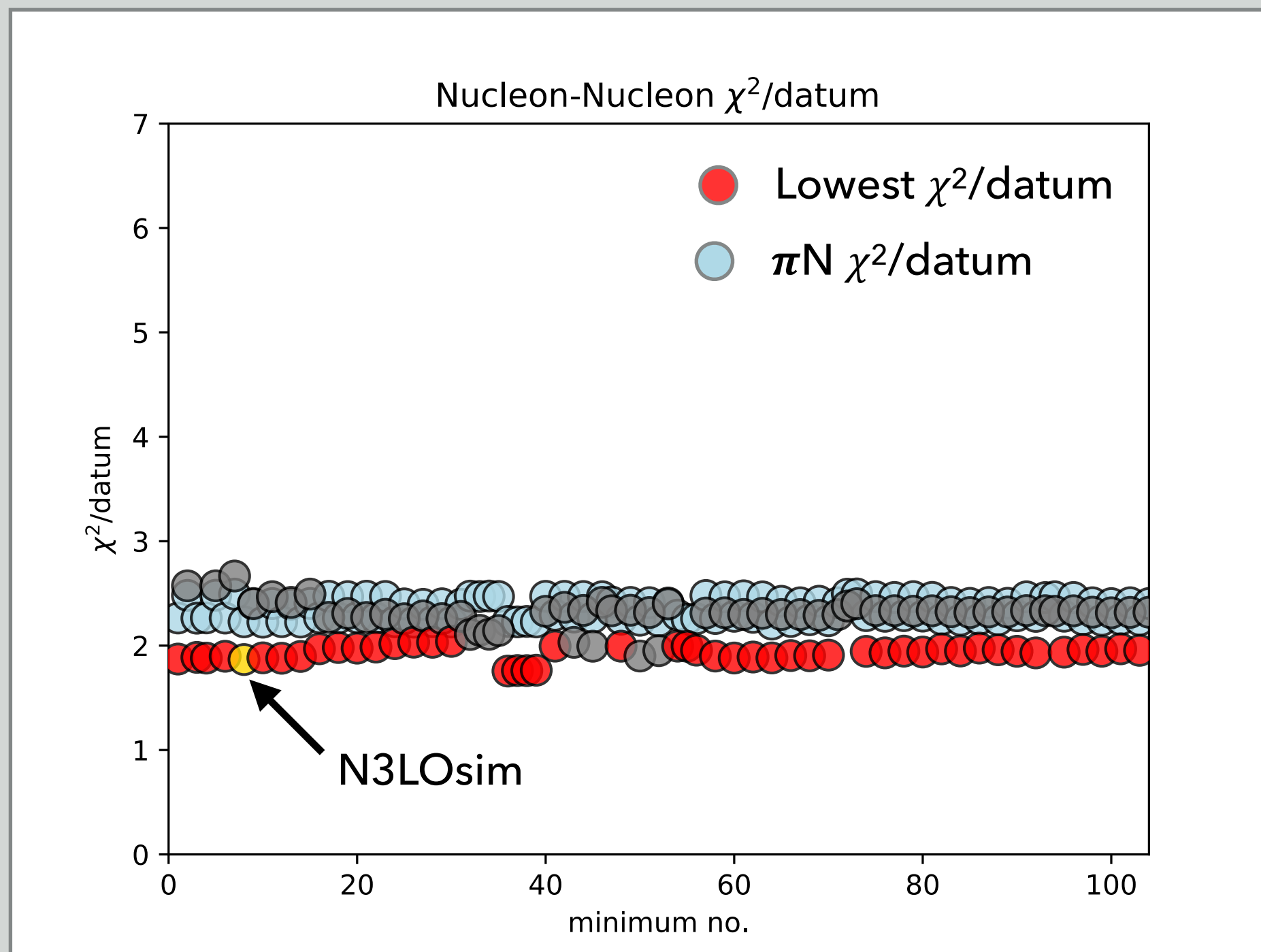
FIT  $C_D, C_E$   
TO  $A=3$  DATA

SIMULTANEOUS  
OPTIMIZATION









FIT  $\pi$ N-SECTOR  
TO  $\pi$ N DATA

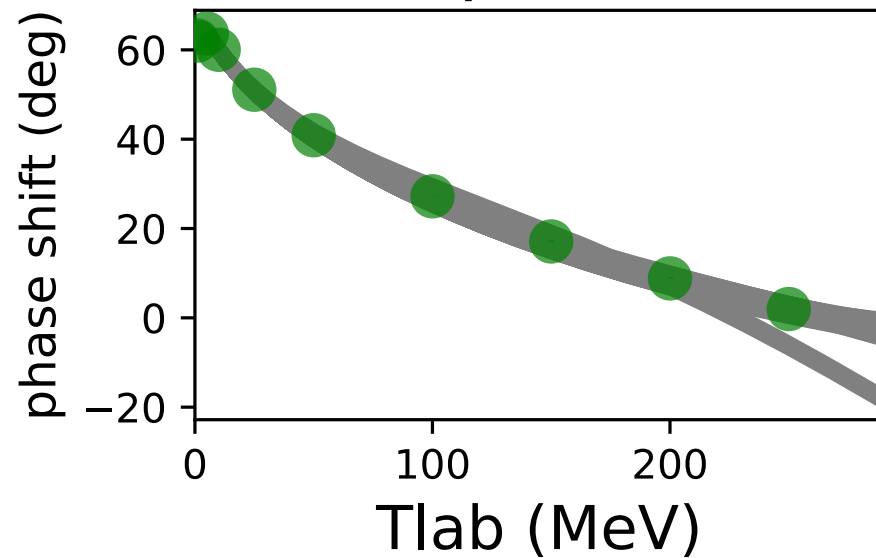
FIT NN  
CONTACTS TO  
NN PHASE SHIFTS

FIT NN  
CONTACTS TO  
NN DATA

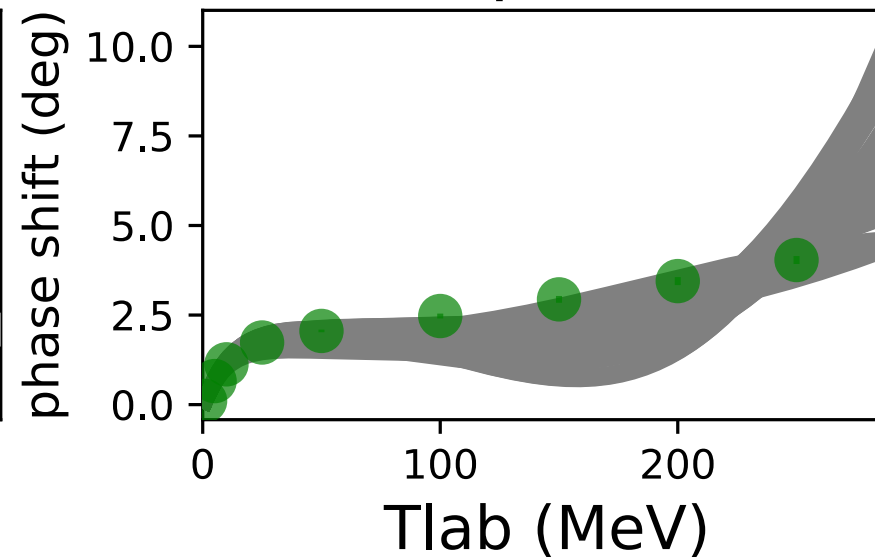
FIT  $C_D, C_E$   
TO  $A=3$  DATA

SIMULTANEOUS  
OPTIMIZATION

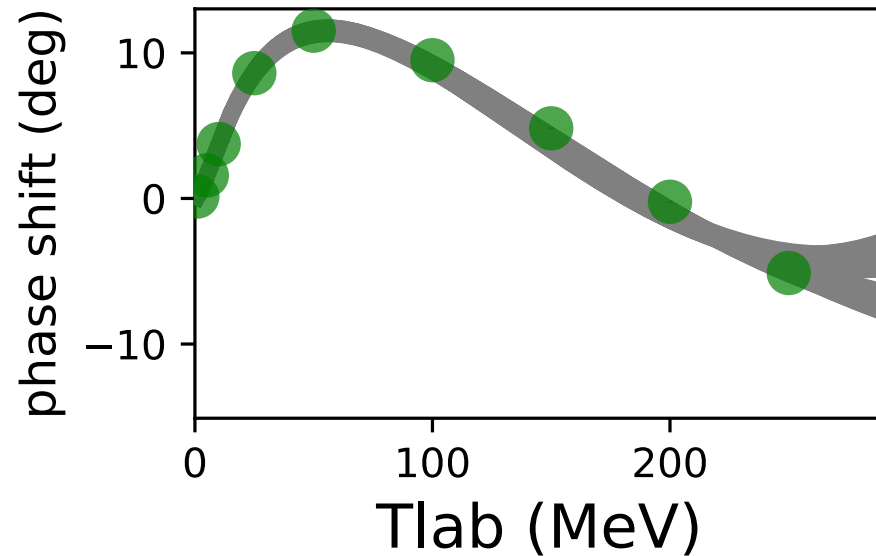
np 1S0



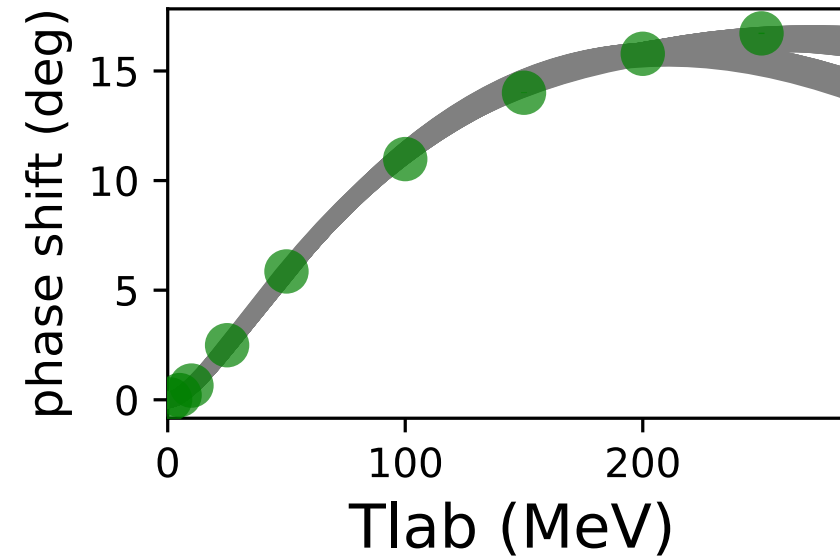
np E1



pp 3P0



pp 3P2



# SCATTERING OBSERVABLES, ORDER-BY-ORDER (SIM)

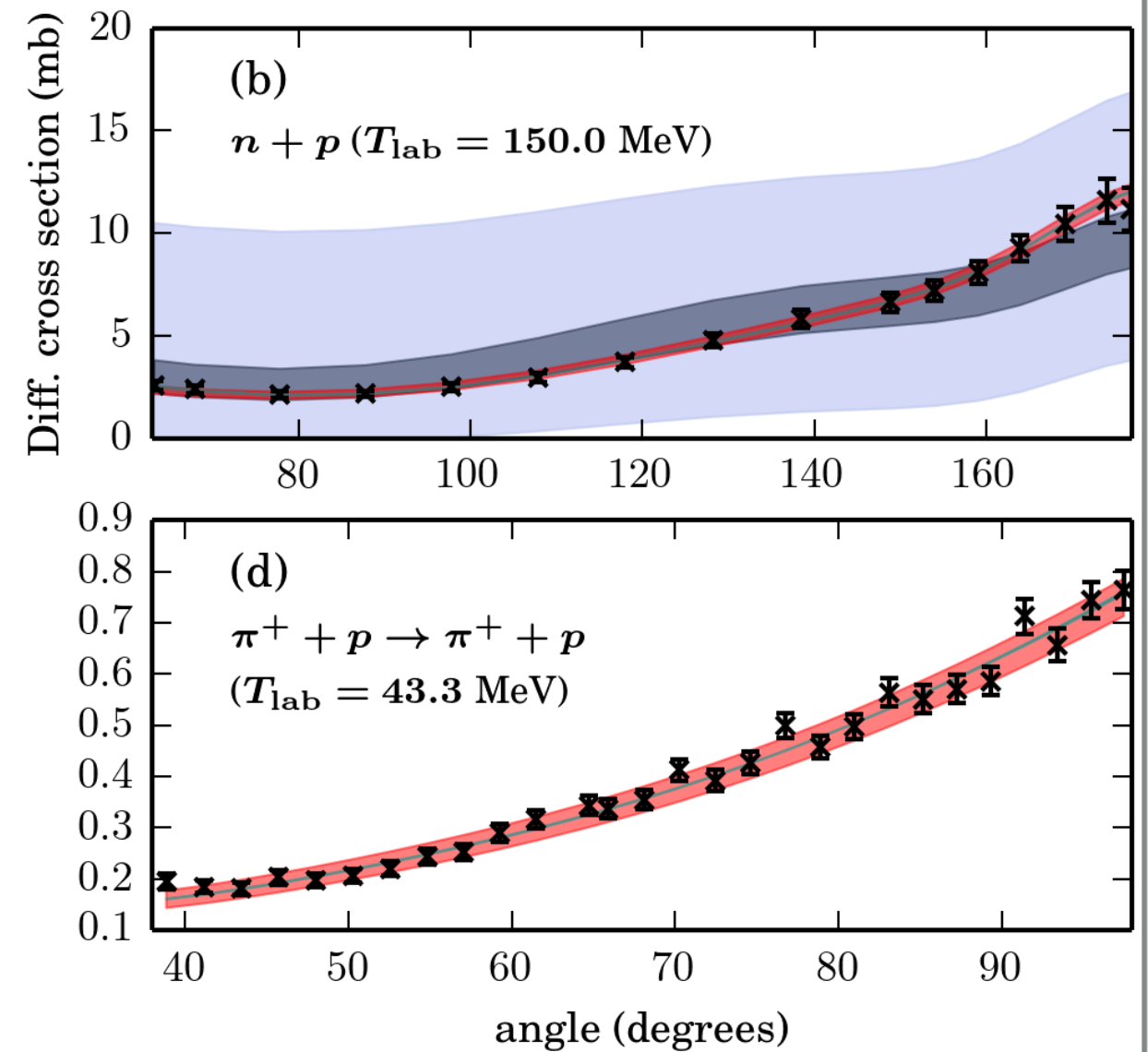
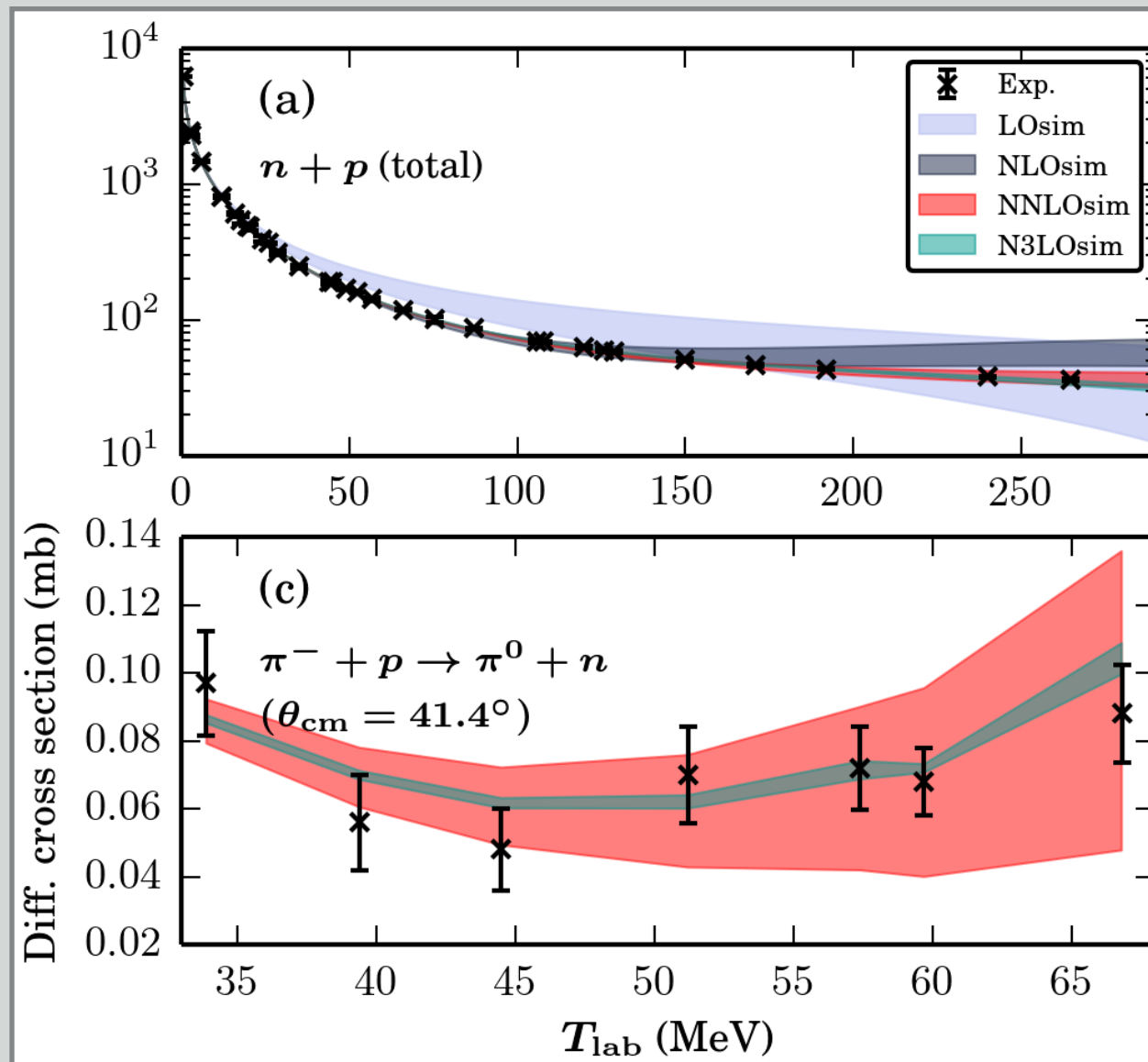
FIT  $\pi N$ -SECTOR  
TO  $\pi N$  DATA

FIT NN  
CONTACTS TO  
NN PHASE SHIFTS

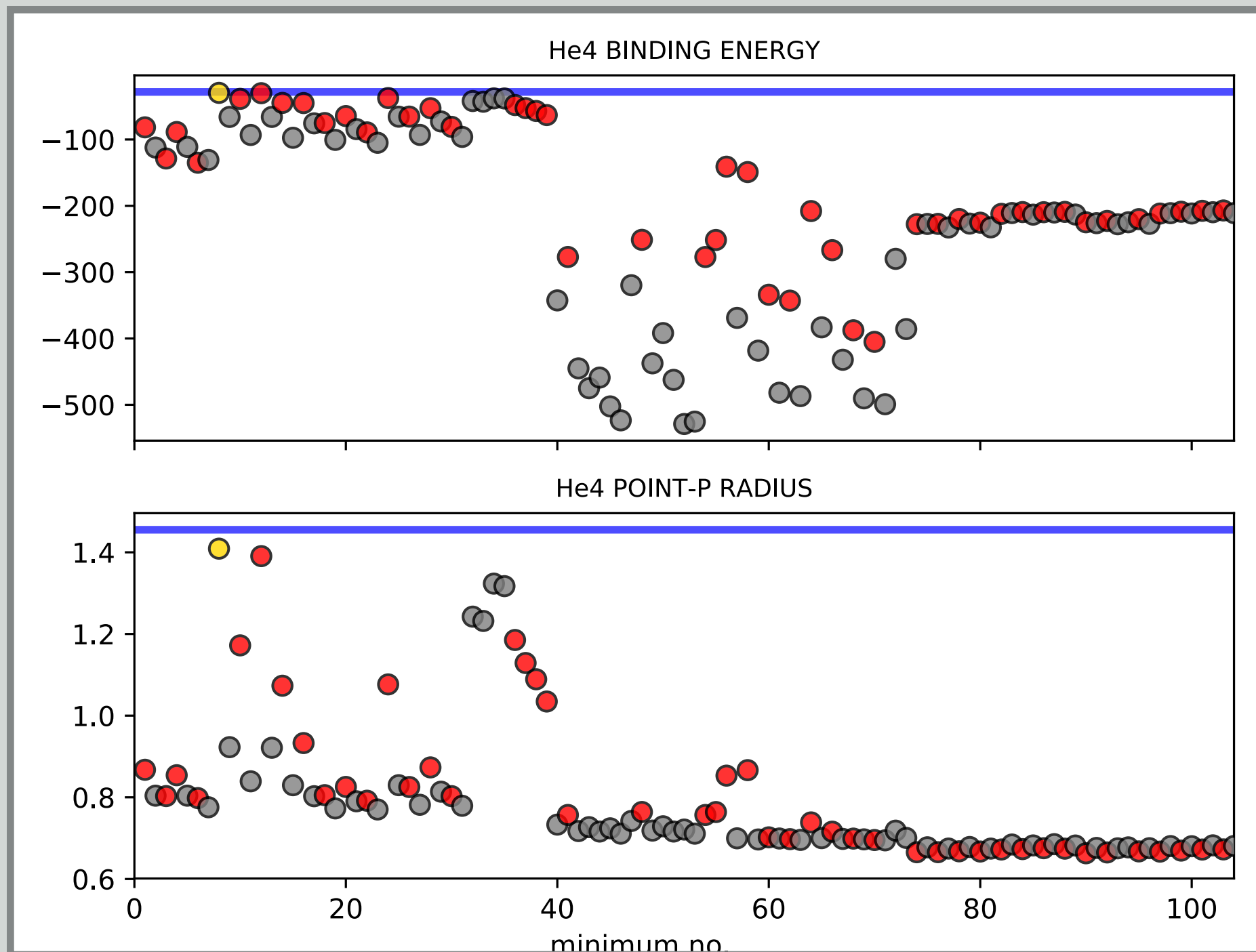
FIT NN  
CONTACTS TO  
NN DATA

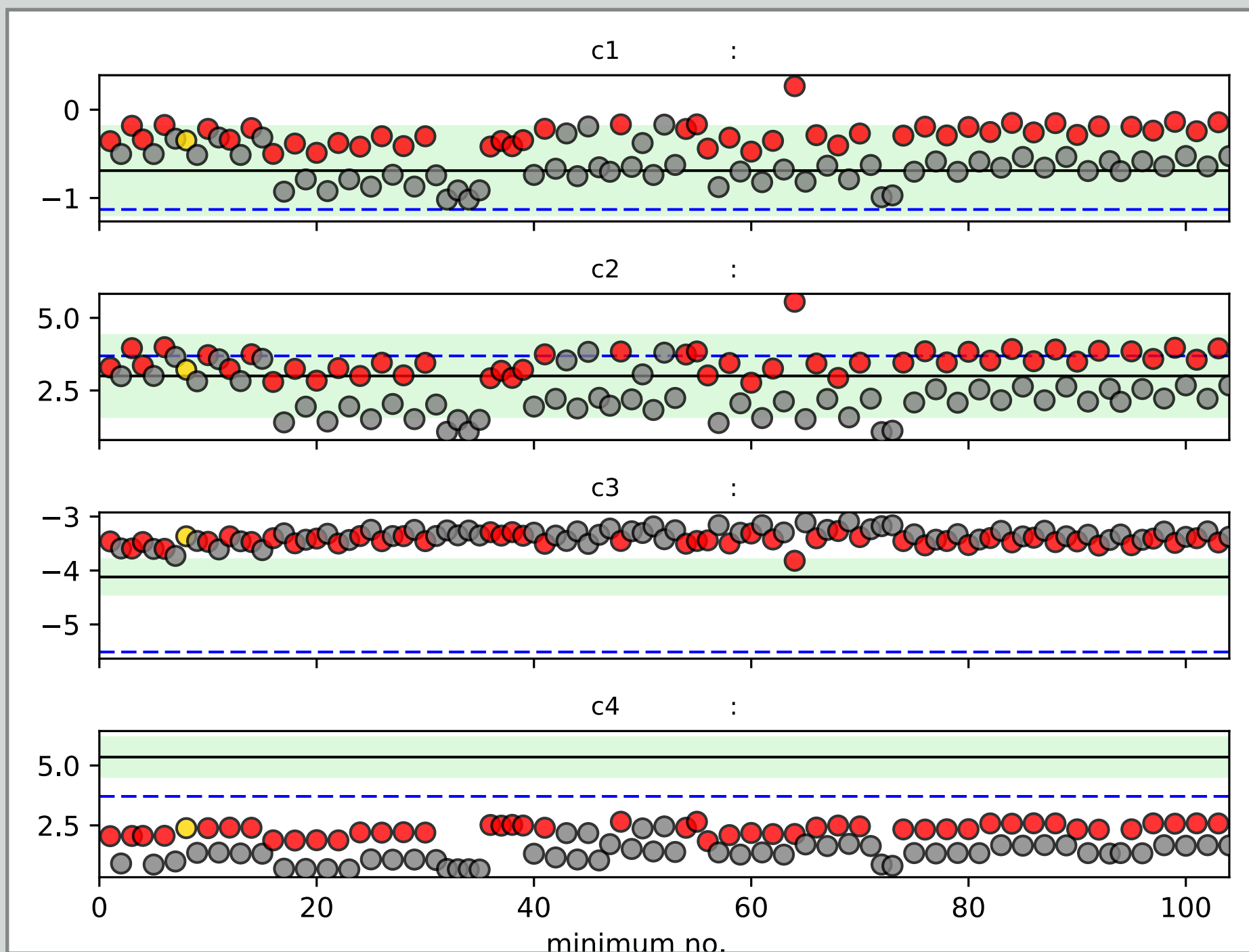
FIT  $C_D, C_E$   
TO  $A=3$  DATA

SIMULTANEOUS  
OPTIMIZATION



# FIRST PREDICTIONS (A=4)





stage 1:  
 $\pi$ N only fit

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 Krebs et al  
 PRC85, 054006



# ROY-STEINER ANALYSIS OF $\pi N$ SCATTERING



- ▶ Physical values of the momentum transfer in  $NN$  scattering is much closer to subthreshold kinematics in  $\pi N$  scattering than to the physical region.
- ▶ Hoferichter et al. recently matched subthreshold parameters of  $\pi N$  scattering from a solution of Roy-Steiner equations to  $\chi$ PT.
- ▶ We allow these results to determine the long-range dynamics of the nuclear force within our optimization framework [ $\pi N$  LECs and covariance matrix from PRL 115, 192301 (2015)].

# ROY-STEINER ANALYSIS OF $\pi N$ SCATTERING



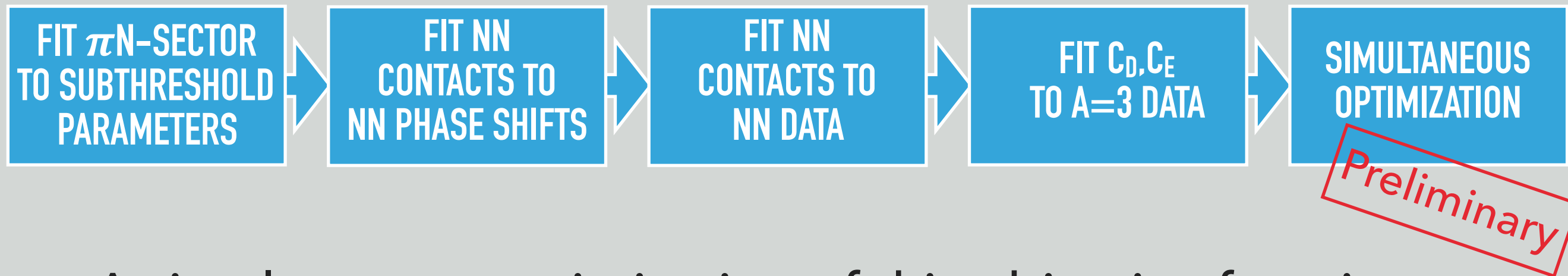
- ▶ We still find multiple minima in the  $NN$  optimization.
- ▶ Keeping only the most promising one we perform the simultaneous optimization of the final stage replacing the optimization w.r.t.  $\pi N$  data with

$$2f(\vec{p}_{\pi N} - \vec{p}_{RS})^T C_{RS}^{-1} (\vec{p}_{\pi N} - \vec{p}_{RS})$$

- ▶ We use  $f=1000$  to stay close to the RS values for the  $c_i$ s.

**Question:** How to best combine experimental data with information from a theory analysis?

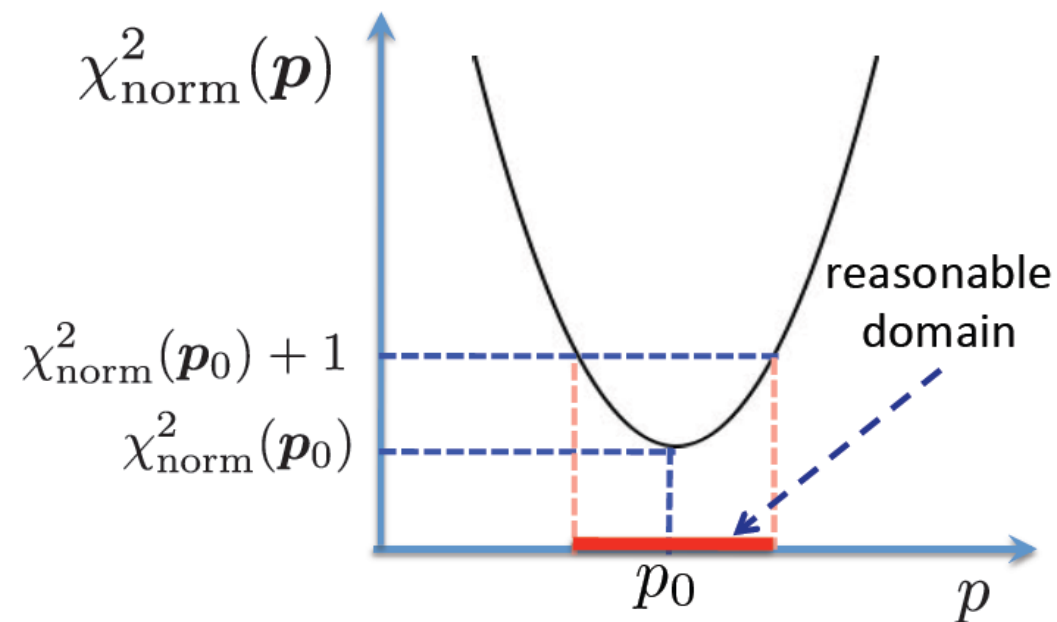
# ROY-STEINER ANALYSIS OF $\pi N$ SCATTERING



- ▶ A simultaneous optimization of this objective function leads to a good description on all  $NN$  and  $NNN$  data.
  - ▶ For instance  $np \chi^2/\text{datum} = 3.5$ .
- ▶ The LECs remain in the Roy-Steiner region (to  $\sim 1\sigma$ )
- ▶ Predictions for  $4\text{He}$  still disagree with experiments
  - ▶  $E(4\text{He}) = -29.4 \text{ MeV}$
  - ▶  $R(4\text{He}) = 1.38 \text{ fm}$

# STATISTICAL ERROR ANALYSIS

- ▶ In a minimum there will be an **uncertainty in the optimal parameter values  $\mathbf{p}_0$**  given by the  $\chi^2$  surface.<sup>1</sup>

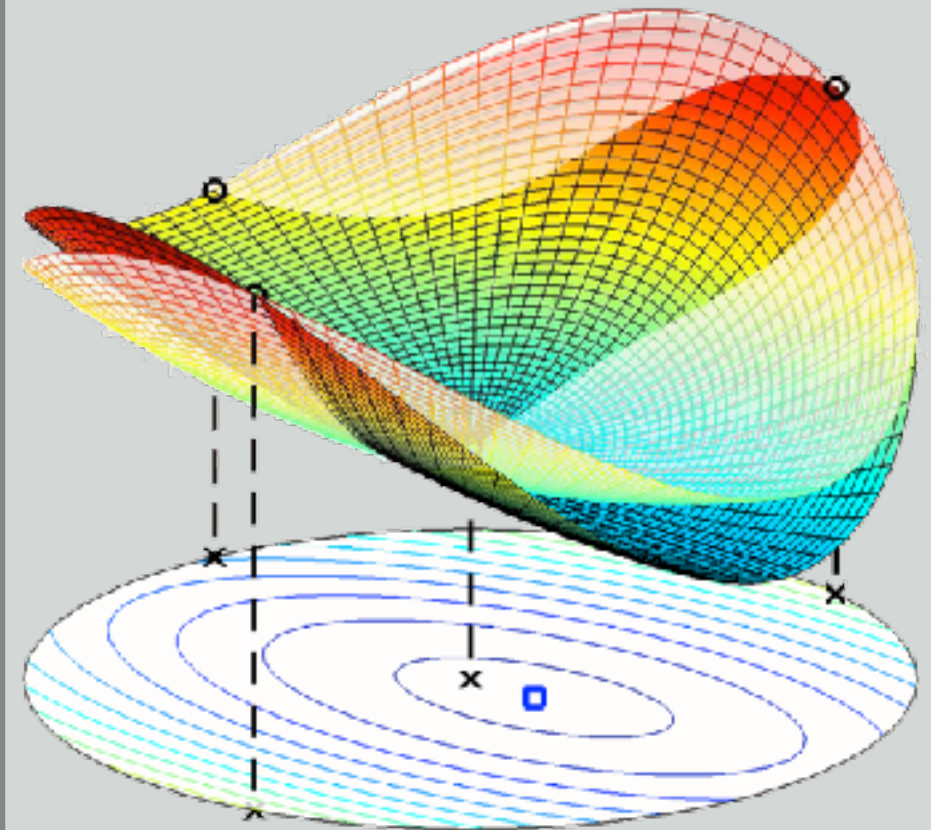


- ▶ Approximate the objective function with a quadratic form in the vicinity of the optimum. Compute the hessian matrix.
- ▶ Expand observables similarly, to second order

$$\mathcal{O}(\mathbf{p}_0 + \Delta\mathbf{p}) - \mathcal{O}(\mathbf{p}_0) \approx (\Delta\mathbf{p}^T) \mathbf{J}_{\mathcal{O}} + \frac{1}{2} (\Delta\mathbf{p}^T) \mathbf{H}_{\mathcal{O}} (\Delta\mathbf{p})$$

- ▶ The covariance between two observables is then

$$\text{Cov}(\mathcal{O}_A, \mathcal{O}_B) \approx \mathbf{J}_{\mathcal{O}_A}^T \text{Cov}(\mathbf{p}_0) \mathbf{J}_{\mathcal{O}_B} + \text{second order}$$



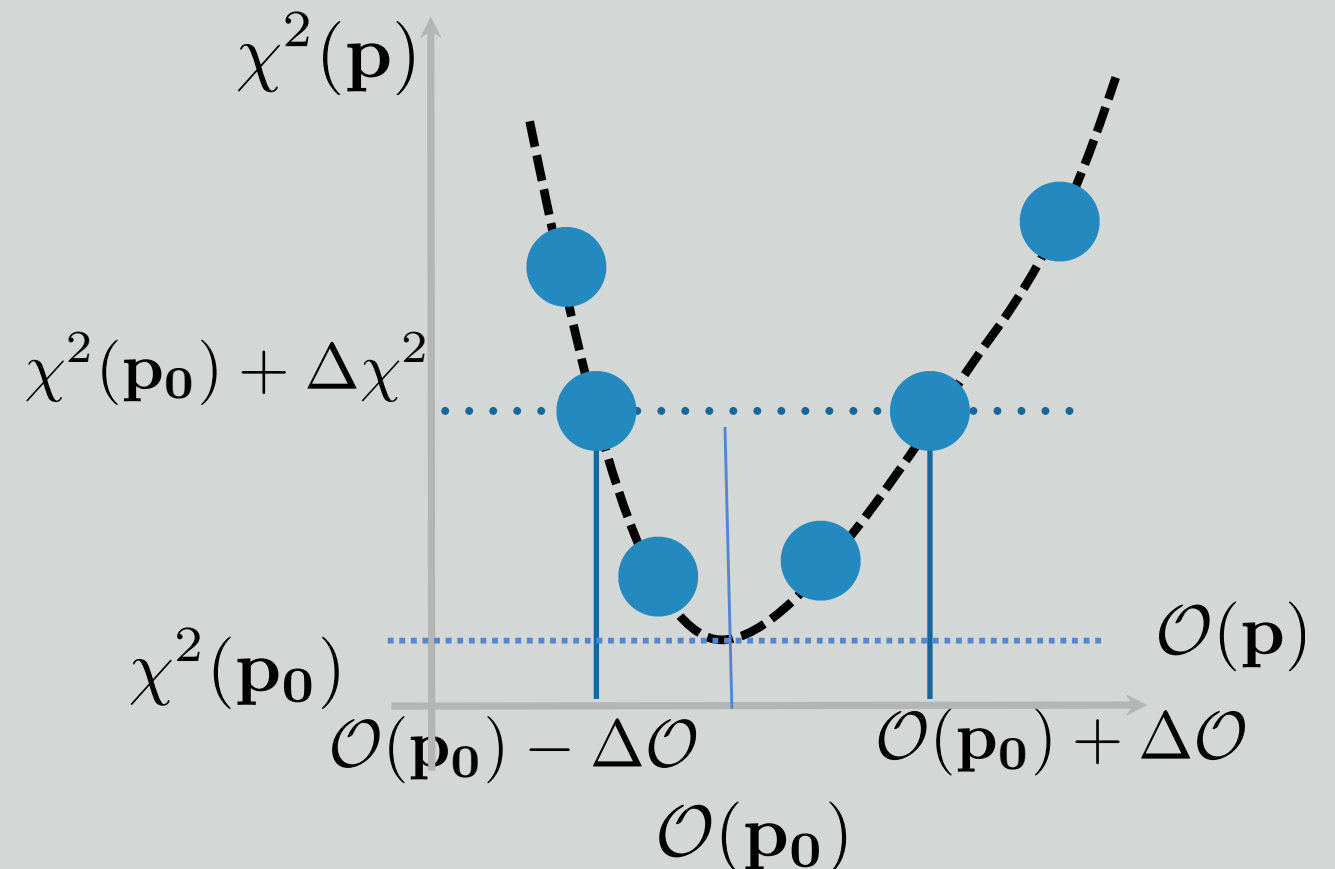
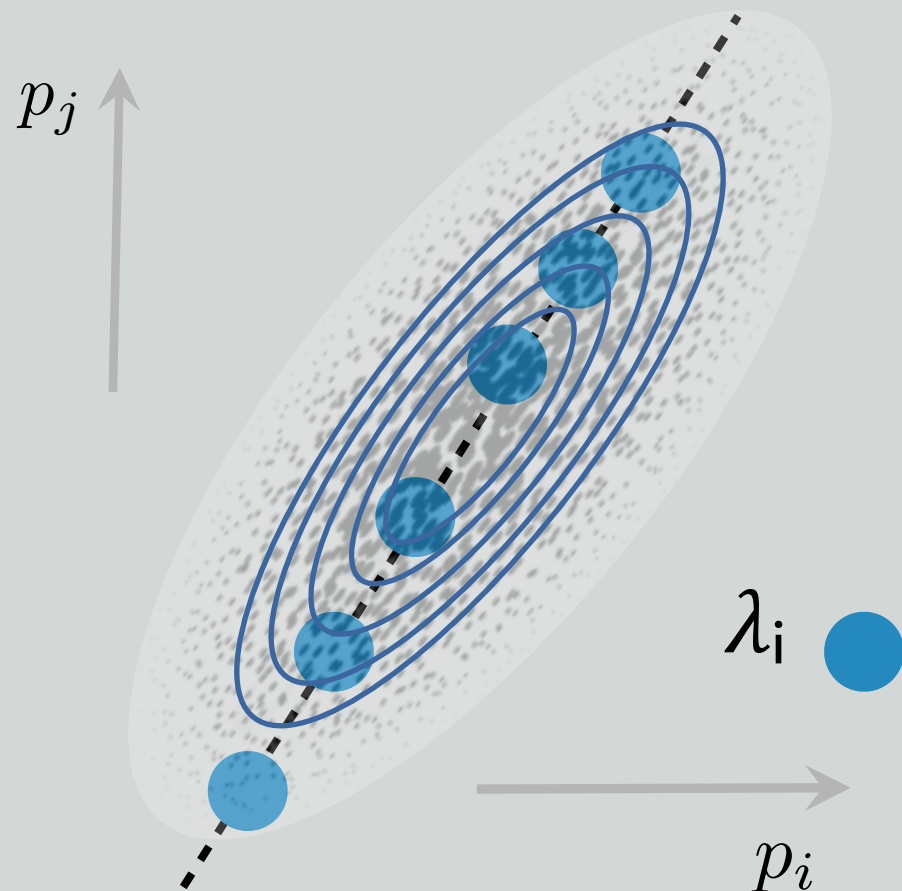
# HESSIAN UNCERTAINTY QUANTIFICATION

	LOsim	NLOsim	NNLOsim	N3LOsim	Exp.
$E(^2\text{H})$	-2.223	$-2.224^{(+1)}_{(-6)}$	$-2.224^{(+0)}_{(-1)}$	$-2.225^{(+1)}_{(-2)}$	-2.225
$E(^3\text{H})$	-11.43	$-8.268^{(+26)}_{(-38)}$	$-8.482^{(+26)}_{(-30)}$	-8.482(28)	-8.482(28)
$E(^3\text{He})$	-10.43	$-7.528^{(+20)}_{(-31)}$	$-7.717^{(+17)}_{(-21)}$	-7.717(19)	-7.718(19)
$E(^4\text{He})$	-40.38(1)	$-27.44^{(+13)}_{(-15)}$	$-28.24^{(+9)}_{(-11)}$	-29.75(18)	-28.30(11)
$r_{\text{pt-p}}(^2\text{H})$	+1.912	$+1.972^{(+0)}_{(-2)}$	$+1.966^{(+0)}_{(-1)}$	+1.977(1)	+1.976(1)
$r_{\text{pt-p}}(^3\text{H})$	+1.292	$+1.614^{(+2)}_{(-3)}$	+1.581(2)	+1.590(2)	+1.587(41)
$r_{\text{pt-p}}(^3\text{He})$	+1.368	+1.791(3)	+1.761(2)	+1.760(2)	+1.766(13)
$r_{\text{pt-p}}(^4\text{He})$	+1.080	+1.482(3)	+1.445(3)	+1.407(5)	+1.455(7)
$E_A^1(^3\text{H})$	-	-	+0.6848(11)	+0.6848(11)	+0.6848(11)
$D(^2\text{H})$	+7.807	$+2.876^{(+85)}_{(-82)}$	$+3.381^{(+46)}_{(-45)}$	$+8.682^{(+96)}_{(-99)}$	-
$Q(^2\text{H})$	+0.3030	$+0.2589^{(+17)}_{(-19)}$	+0.2623(8)	+0.2897(7)	+0.270(11)
$a_{nn}^N$	-26.04(8)	$-18.95^{(+38)}_{(-41)}$	$-19.28^{(+74)}_{(-80)}$	$-19.50^{(+79)}_{(-89)}$	-18.95(40)
$a_{np}^N$	-25.58(8)	$-23.60^{(+10)}_{(-13)}$	-23.83(11)	$-23.75^{(+3)}_{(-6)}$	-23.71
$a_{pp}^C$	-7.579(6)	$-7.799^{(+1)}_{(-3)}$	-7.811(1)	$-7.812^{(+2)}_{(-6)}$	-7.820(3)
$r_{nn}^N$	+1.697(1)	$+2.752^{(+7)}_{(-8)}$	+2.793(14)	+2.785(15)	+2.75(11)
$r_{np}^N$	+1.700(1)	+2.648(3)	+2.686(2)	+2.683(2)	+2.750(62)
$r_{pp}^C$	+1.812(1)	+2.704(3)	+2.758(2)	+2.755(2)	+2.790(14)

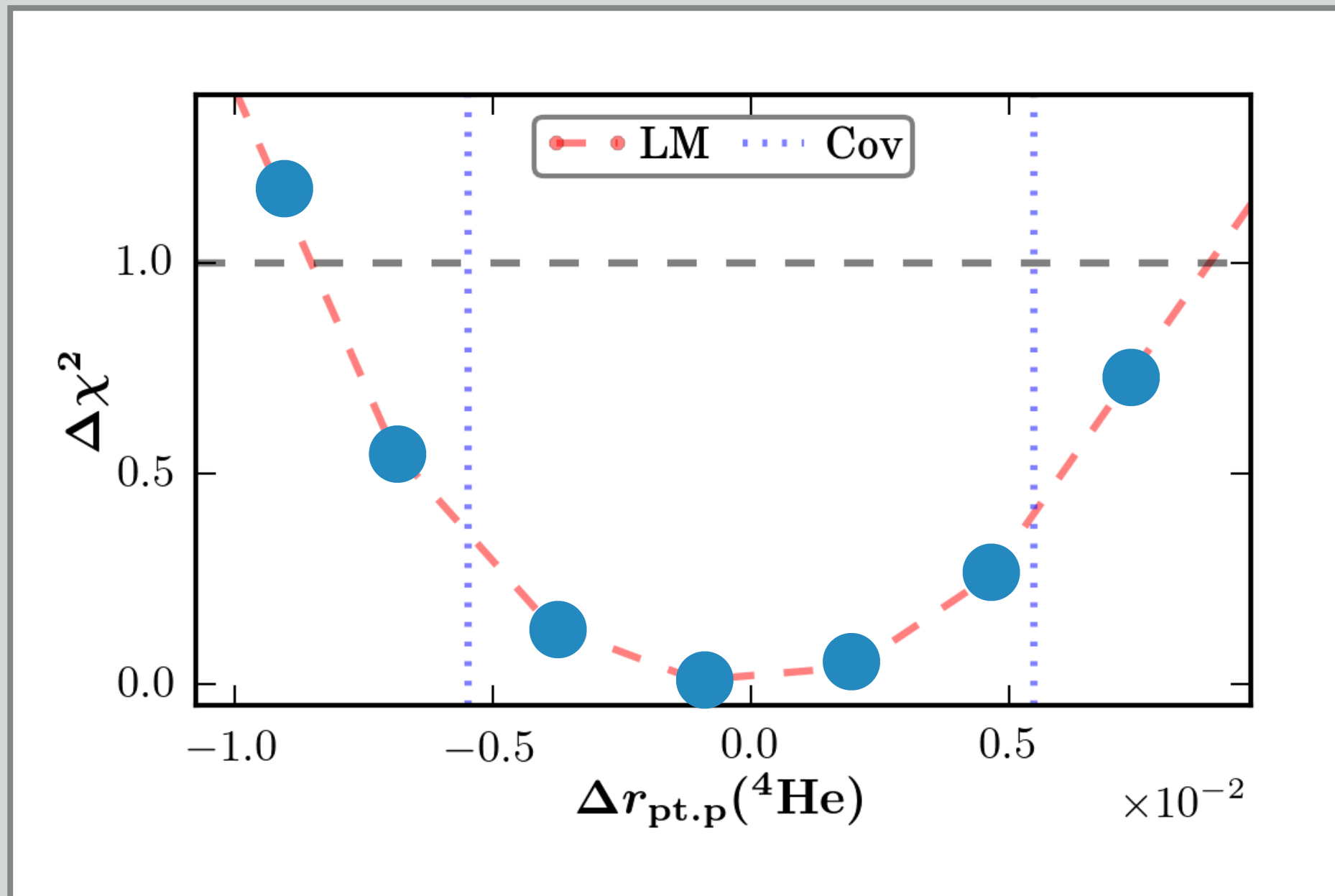
# LAGRANGE MULTIPLIER

- ▶ The hessian UQ relies on a quadratic shape of the  $\chi^2$ -surface.
- ▶ The local minima of our  $\chi^2$ -surface are often very flat in some directions. Curvature is dominated by fourth-order terms.
- ▶ A possible solution: Lagrange multiplier optimization:

- ▶ Minimize  $F(\mathbf{p}, \lambda) = \chi^2(\mathbf{p}) + \lambda \cdot \mathcal{O}(\mathbf{p})$



# COMPARISON: HESSIAN UQ VS LAGRANGE-MULTIPLIER UQ



see also B. Carlsson et al., Phys. Rev. C 95 (2017) 034002



# EXPLORING FURTHER SYSTEMATIC UNCERTAINTIES

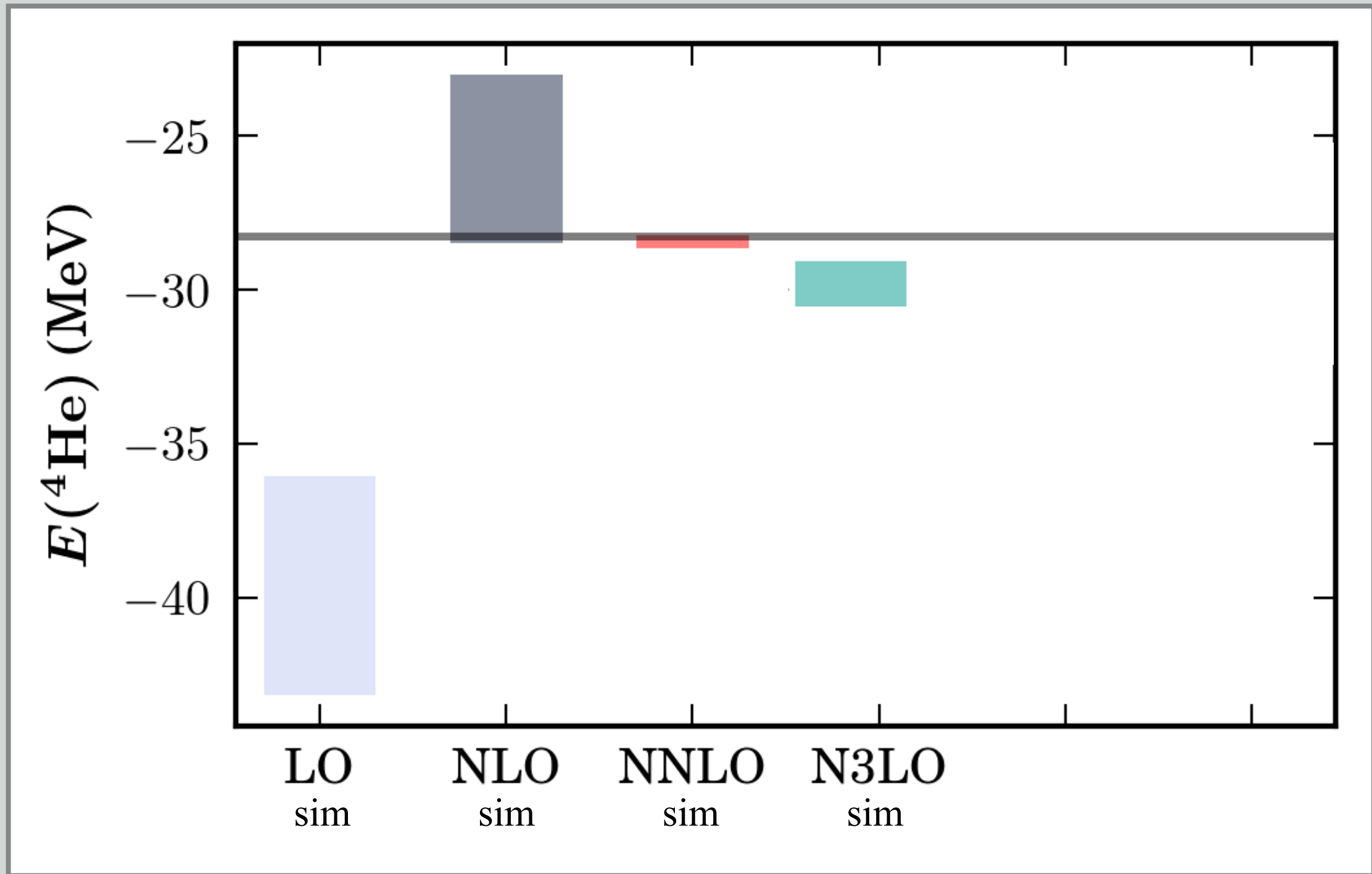
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- ▶ So far, all results have been obtained with a non-local regulator with cutoff  $\Lambda=500$  MeV.
  - ▶ A subset of systematic uncertainties can be probed by varying  $\Lambda$ .
- ▶ Reoptimizing with different  $\Lambda$  (450-575 MeV) will give us a **family** of models.
- ▶ **All of them will reproduce the same few-body physics.**



# SYSTEMATIC UNCERTAINTIES

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**CONCLUSION**

# SUMMARY (MAINLY FOR PHYSICISTS)

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- ▶ There are roughly 100 local minima when the non-local N3LO  $NN+NNN$  interaction is optimized w.r.t.  $A=2,3$  data.
  - Typically, the  $NN$  and  $\pi N$   $\chi^2/\text{datum}$  is 2 across the board.
  - The best few (2-4) candidate(s) predict the  $E(4\text{He})$  within 2 MeV, however the radii are too small.
- ▶ In few-nucleon calculations, the non-local 3N-interaction is not a small perturbation, compared to N2LO,
- ▶  $\pi N$  coupling constants are of expected size.
  - However, they are very poorly constrained from  $\pi N$  data alone.
  - When fitted simultaneously,  $c_3$   $c_4$  deviate significantly from the values obtained when fitted w.r.t.  $\pi N$  scattering data.
- ▶ Lagrangian multiplier optimization and UQ for non-quadratic cases. No need for derivatives.

# OUTLOOK

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- ▶ The inclusion of more data in the objective function requires other approaches to the optimization problem needed. (See also Andreas' presentation.)
- ▶ The frequentist approach does not offer an easy and transparent method for handling systematic uncertainties or imposing prior knowledge.
- ▶ Bayesian parameter estimation is advantageous, but costly.
  - avoiding the need to 'judge', a priori, what data can be included in order to safely avoid overfitting.
  - not obvious whether local minima will vanish.
  - offers a viable approach to include prior knowledge of certain parameters from Roy-Steiner analysis.
- ▶ Investigate other chiral EFT power-counting schemes. (See also Andreas' presentation.)